

# Electron Clouds in the PSR and SNS

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# Collaborators

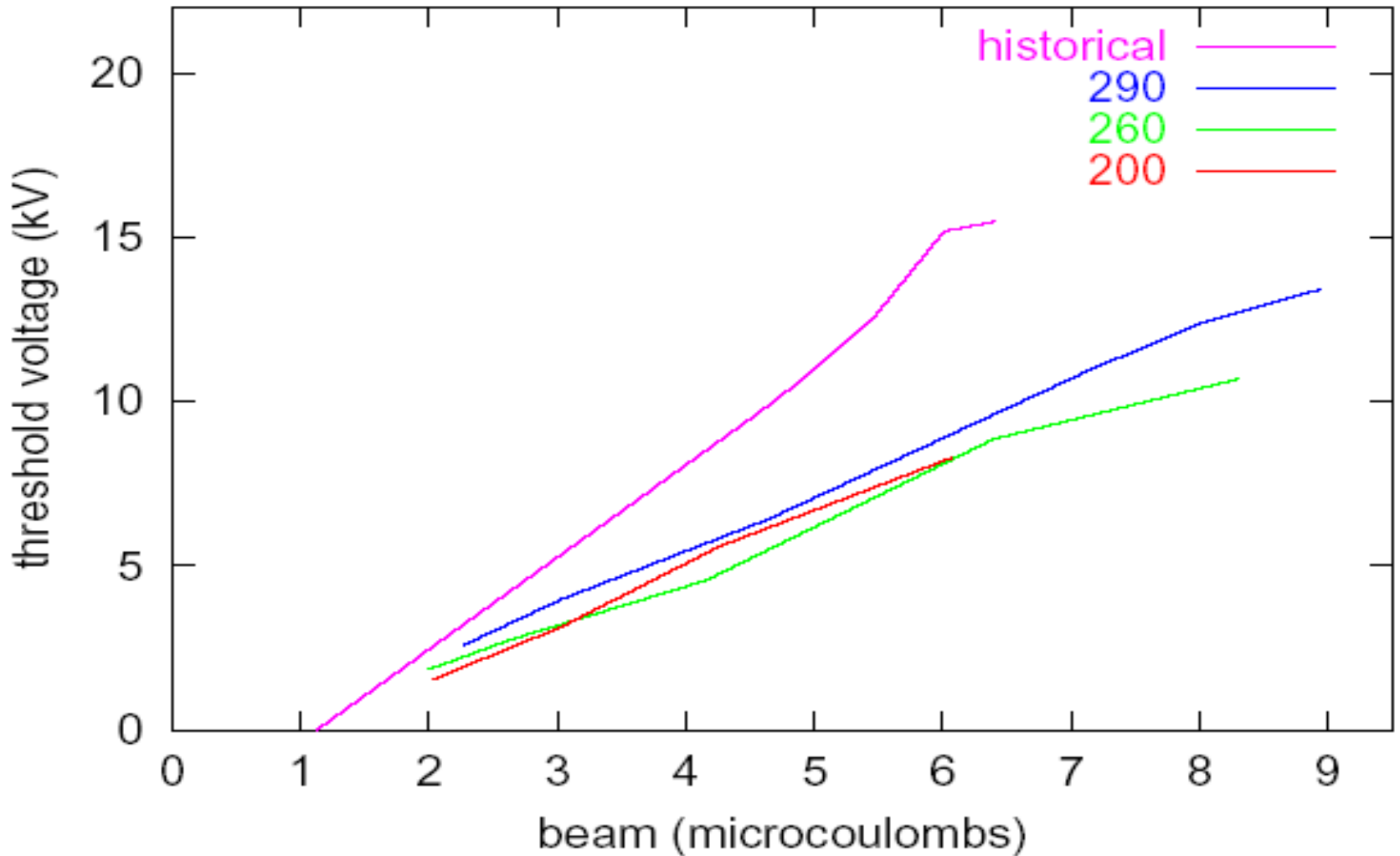
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- LANL: R. Macek
- LBNL: M. Furman, M. Pivi

- **Accelerator physics issues**
  - Electron cloud generation
  - Stability simulation
  - Stability eigenmode analysis

# PSR Threshold before and after long run for a range of bunch lengths (Courtesy R. Macek)



Primary electrons are generated via losses and collisions with gas

Model loss rate proportional to instantaneous current

Along with secondary emission there is also the possibility of reflection

Microphysics gets very complicated

A simple parameterization is given by

$$R = (R_0 - R_\infty) \exp(-E / E_{char}) + R_\infty$$

$$\delta(E) = \delta_{max} \times 1.1 \frac{1 - \exp[-2.3(E / E_{max})^{1.35}]}{(E / E_{max})^{0.35}}$$

If not reflected then use a Lorentzian distribution for secondary energy.

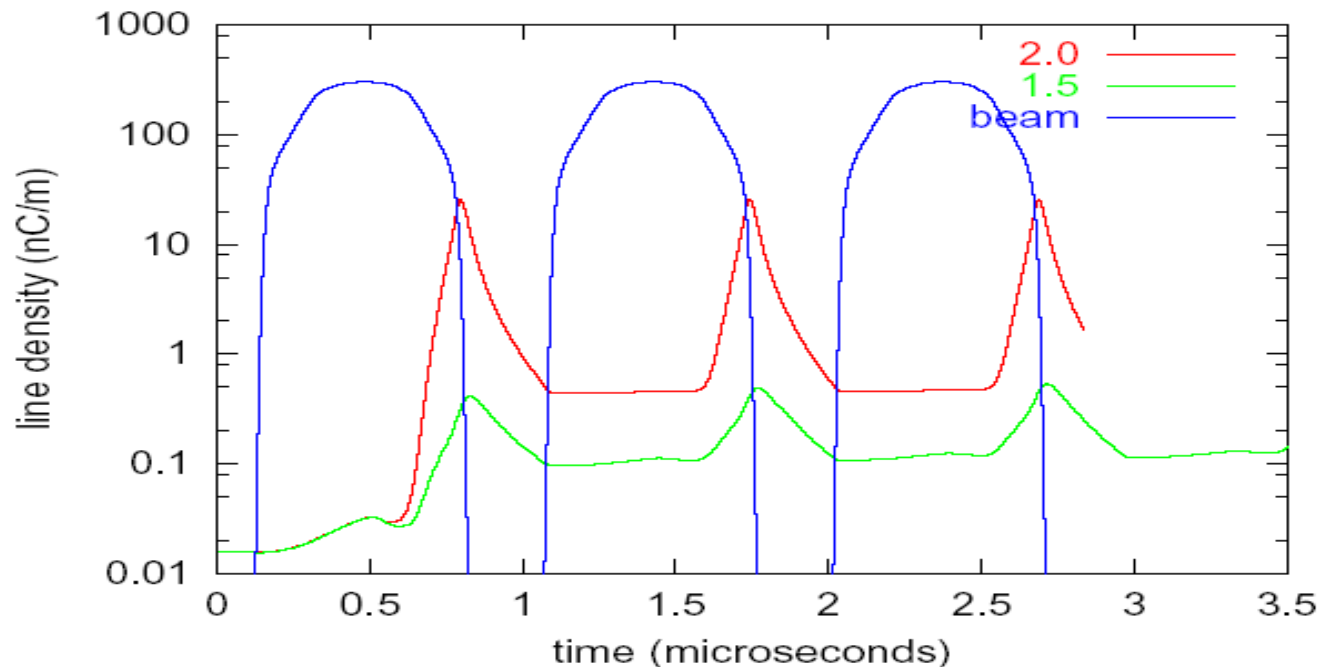
Wall conditions are somewhat uncertain for SNS

electron generation rates and reflectivity are important

Uses a fast cylindrically symmetric algorithm and simple SEY model

2 MW SNS, with  $E_{\max}=300$  eV,  $\delta_{\max}=1.5$  and 2.0, loss rate  $\propto I(t)$

200 e per lost proton and 0.1% losses gives 2E8 e/m/turn



# Electron-Proton Instability coasting beam scaling



- Characteristic tunes

$$Q_e^2 = \frac{e\lambda_p}{2\pi\epsilon_0 a^2 m_e \omega_0^2}, Q_p^2 = \frac{e\lambda_e}{2\pi\epsilon_0 a^2 \gamma m_p \omega_0^2}$$

- Coasting beam tune shift

$$\Delta Q_0 = \Delta Q_{sc} + i \frac{Q_p^2 Q_e}{2Q_\beta \delta Q_e}$$

- Coasting beam tune spread

$$\delta Q_\beta = \left| \eta \right| Q_e \frac{\Delta p}{p} \Big|_{hwhm}$$

- Threshold assuming SC dominates  
magnitude of tune shift

$$N_p \propto V_{rf} \tau_b^3$$

- For production length PSR beams  
the proportionality constant is  
reasonable
- Observed PSR threshold is roughly  
independent of bunch length
- Scaling with bunch length  
and inclusion of realistic clouds  
have been focus of simulations

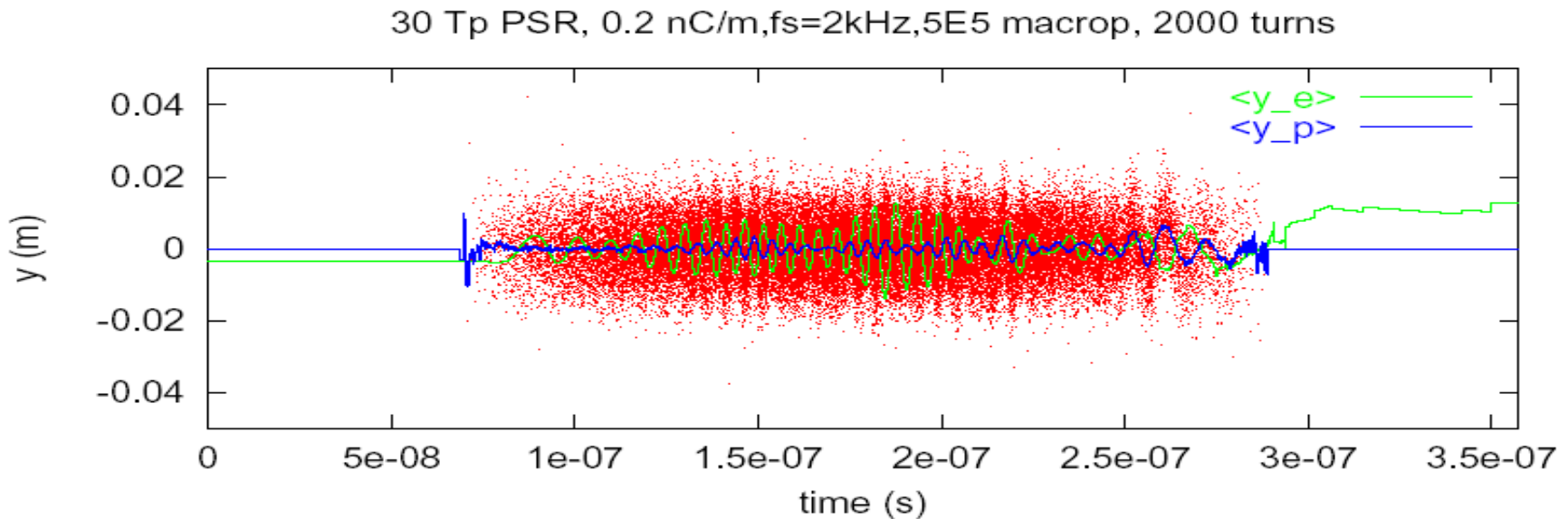
- 20 to 200 electron cloud macroparticles and space charge kick implemented  $\approx 10$  times per betatron oscillation
- Cloud density at head of bunch is an input parameter
- Only vertical coherent forces on protons and linear RF force
- Contribution from straight sections greater than dipoles, model straights
- Electrons can multipactor, no loss driven production
- Macroparticle equations of motion

$$\frac{d^2 y_p}{d\theta^2} = -Q_y^2 y_p + C_{sc} \lambda_p(\theta, t) (y_p - \langle y_p(\theta, t) \rangle) + F_e(\theta, y_p, t)$$

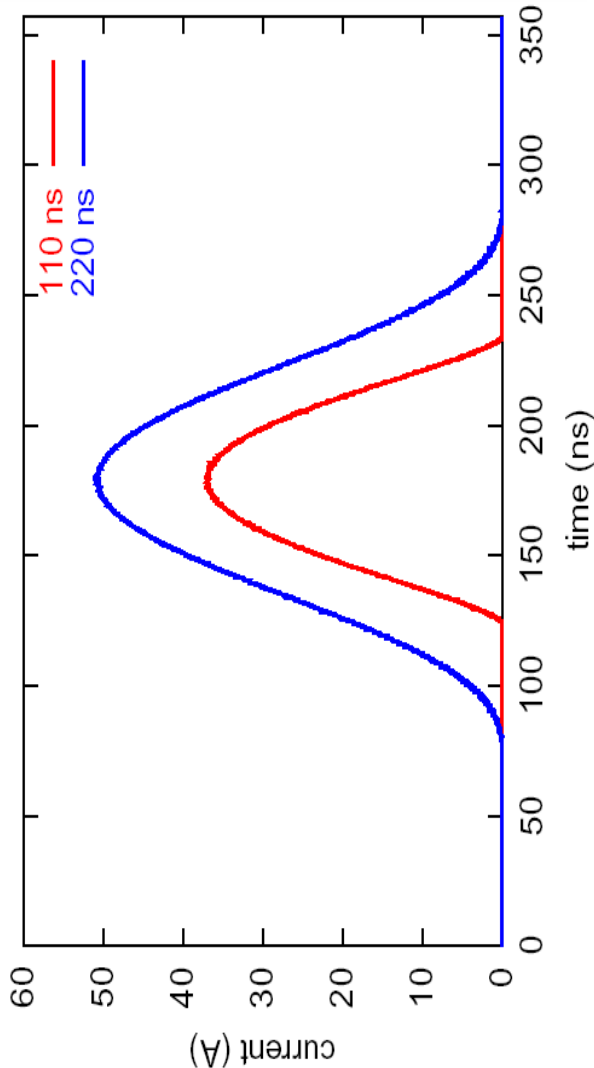
$$F_e = - \frac{e \lambda_e(\theta, t)}{2\pi\epsilon_0} \frac{y_p - \langle y_e(\theta, t) \rangle}{\sigma^2(y_e(\theta, t)) + \Delta x^2 + \Delta y^2} \frac{1}{\gamma m_p \omega_0^2}$$

$$m_e \frac{d^2 y_e}{dt^2} = - \frac{e \lambda_p(\theta, t)}{2\pi\epsilon_0} \frac{y_e - \langle y_p(\theta, t) \rangle}{\max(a^2, \Delta x^2 + \Delta y^2)}$$

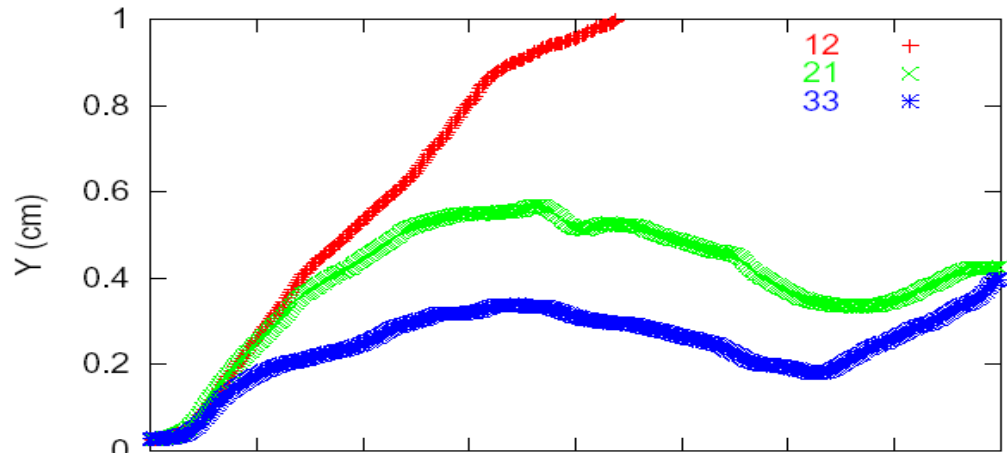
- Did simulations for a range of bunch lengths, intensities, RF voltage, and electron densities.
- Electron densities of 1nC/m generally led to an explosive instability when other parameters were well below thresholds observed in the PSR.
- For smaller electron densities, found that beams that were unstable at small amplitude could reach a saturation amplitude and continue oscillating without filamentation.
- Data taken by R. Macek in PSR suggest a similar behavior in real beams with saturation amplitudes of order a millimeter



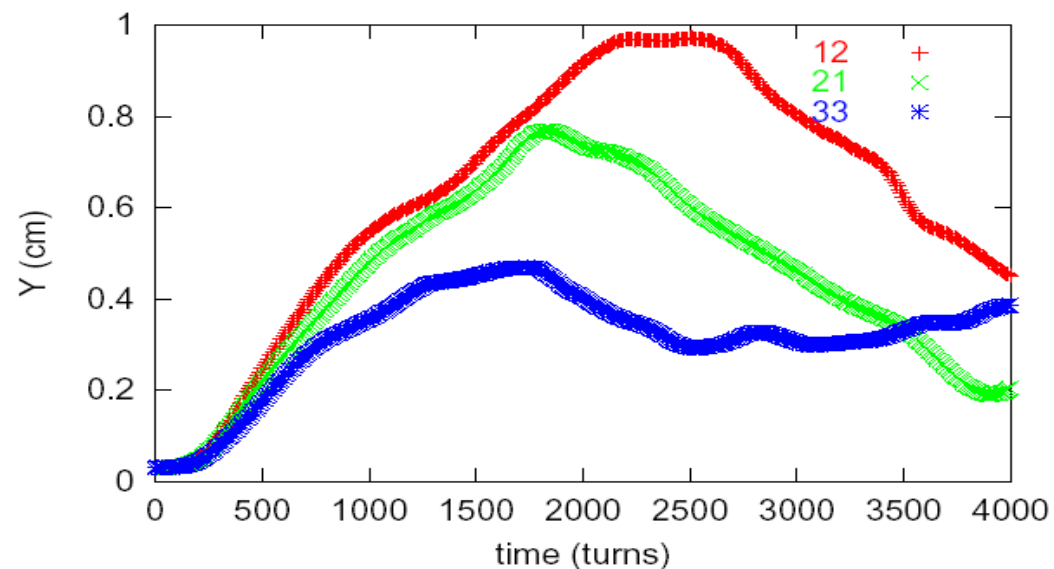
# PSR simulations $Y^2 = \int (\langle y(t) \rangle^2 + \langle y'(t)/Q_y \rangle^2) \rho(t) dt$



PSR, 220 ns, 4.8 uC, 0.2 nC/m, various Vrf (kV)



PSR, 110 ns, 2.4 uC, 0.2 nC, various Vrf (kV)

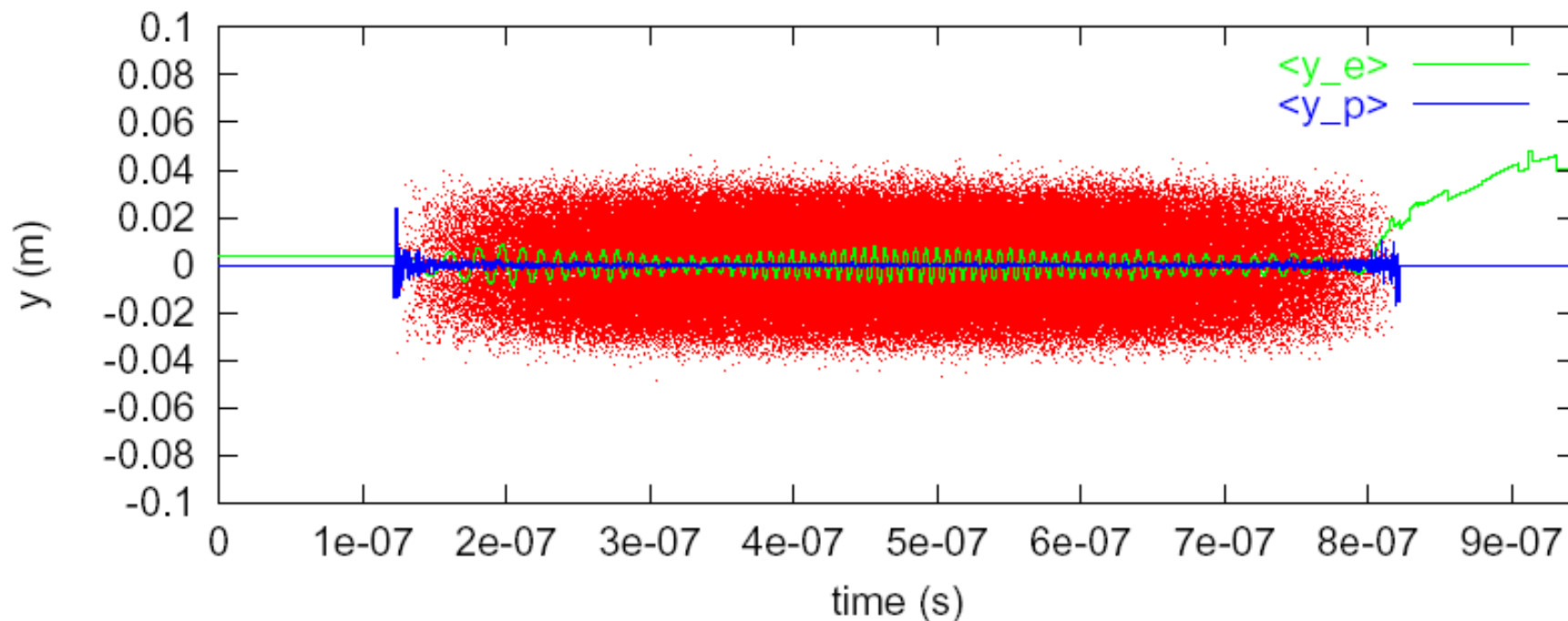


# Comments on PSR simulations and SNS results



- Simulations give threshold electron densities 20% of observed
- Simulations need about 50% more RF than observed (11 $\mu$ H inductor)
- Unclear where additional damping comes from
- SNS predictions should be conservative

2MW SNS 1 nC/m, 2E6 macrop, 2000 turns



$$\frac{\partial x(\varphi, \nu, \theta)}{\partial \theta} + \nu \frac{\partial x}{\partial \varphi} + F_{\text{well}}(\varphi) \frac{\partial x}{\partial \nu} = i\Delta Q_{sc} [x(\varphi, \nu, \theta) - \bar{x}(\varphi, \theta)]$$

$$+ i \frac{Q_p^2 Q_e^2}{2Q_x \tilde{Q}} \int_0^\varphi \bar{x}(\varphi_1, \mathcal{G}) \sin[\tilde{Q}(\varphi - \varphi_1)] \exp[-Q_e(\varphi - \varphi_1) / 2Q_r] d\varphi_1$$

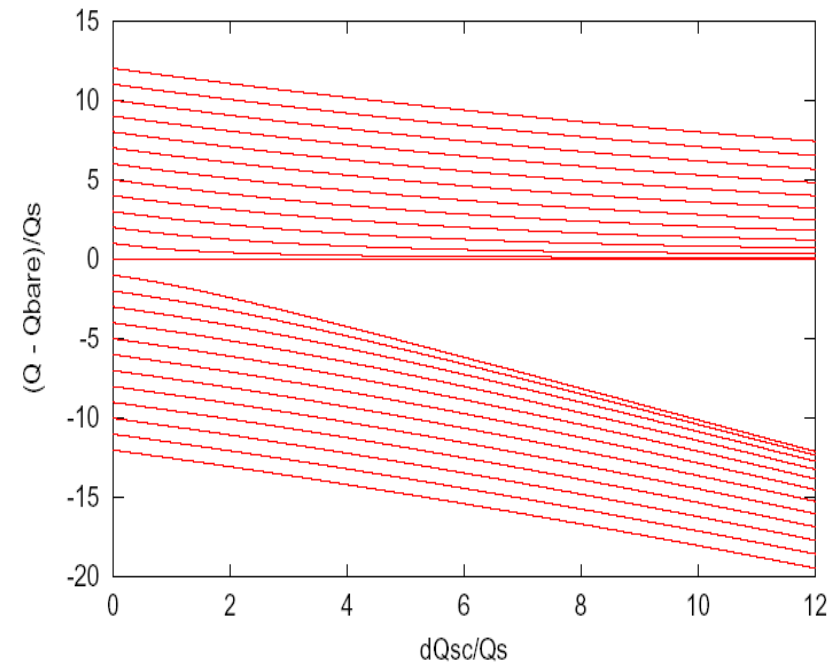
Approximate RF by a square well potential

Need tune eigenvalues of the PDE

This is fairly involved

Details in upcoming article

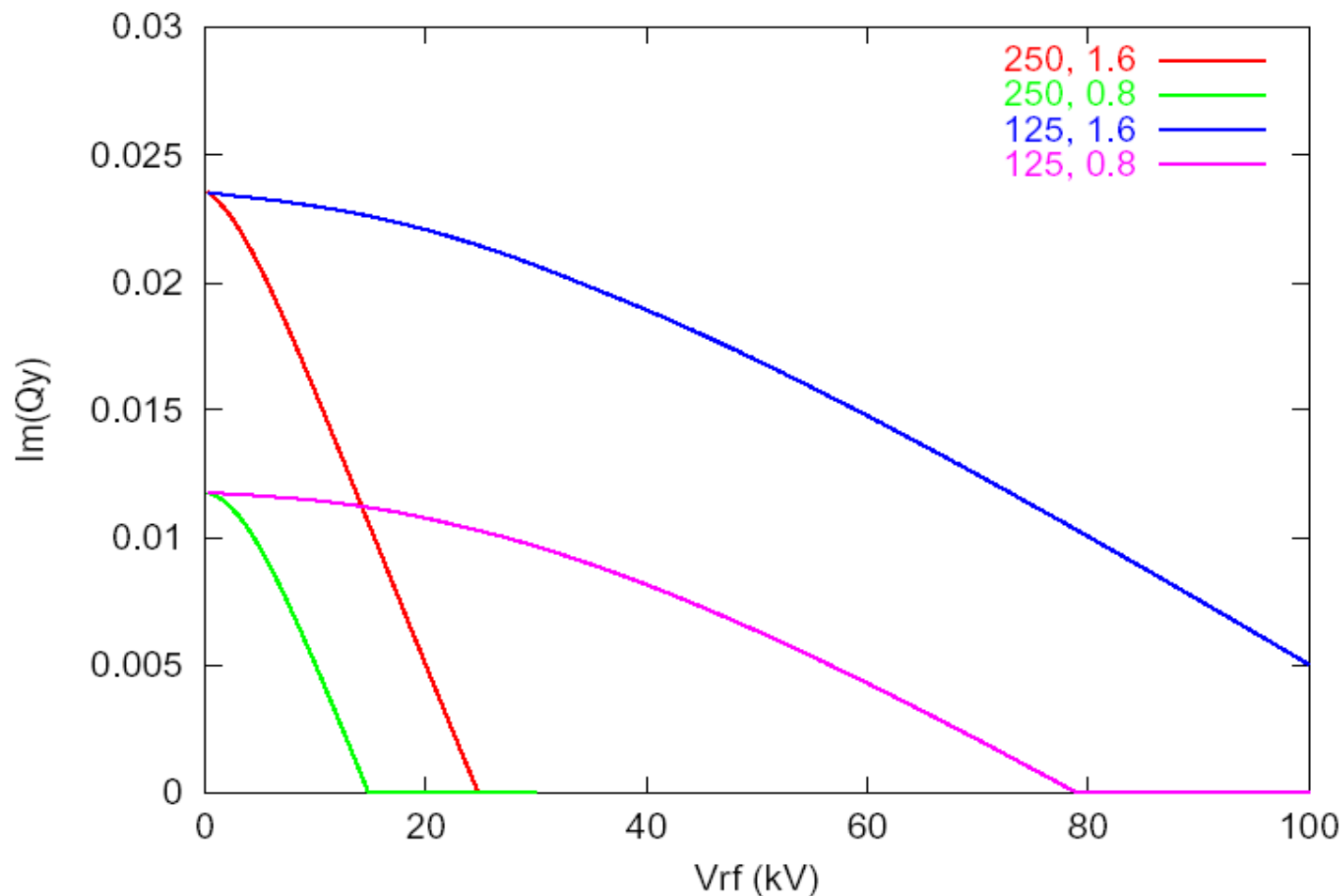
$$\Delta Q_{sc} \approx 100Q_s$$



# Growth rates versus rf voltage

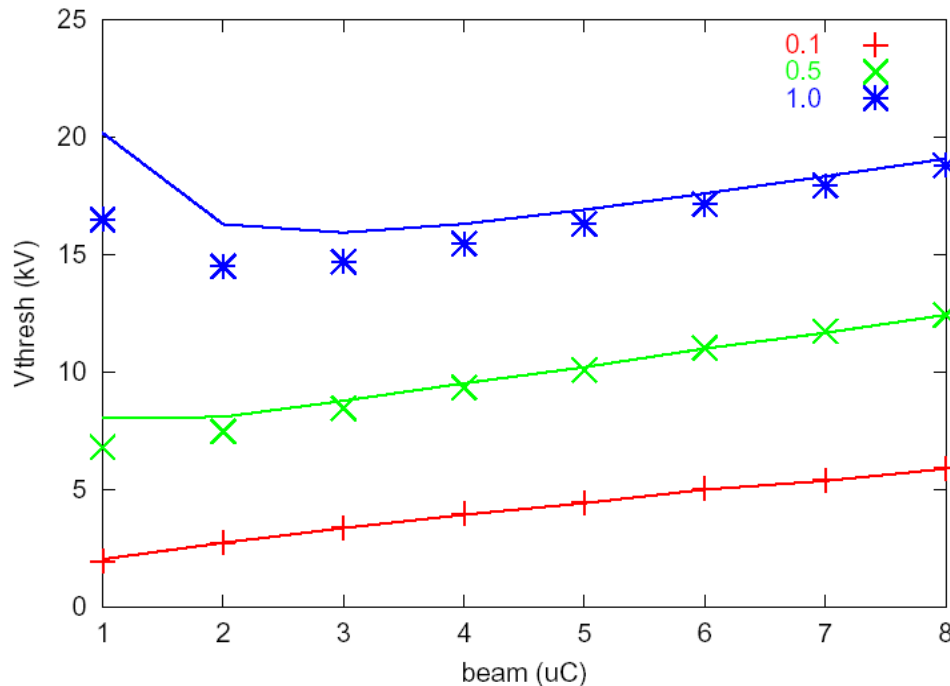
## PSR growth rates

For various bunch lengths (ns) and electron densities (nC/m)



# Coasting beam versus eigenmodes

The symbols are the threshold from the eigenmode analysis for various electron densities (pure TMCI)  
The solid line is the coasting beam threshold with same tune distribution



$$\Delta Q_0 = \Delta Q_{sc} + i \frac{Q_p^2 Q_e}{2Q_\beta \delta Q_e}$$

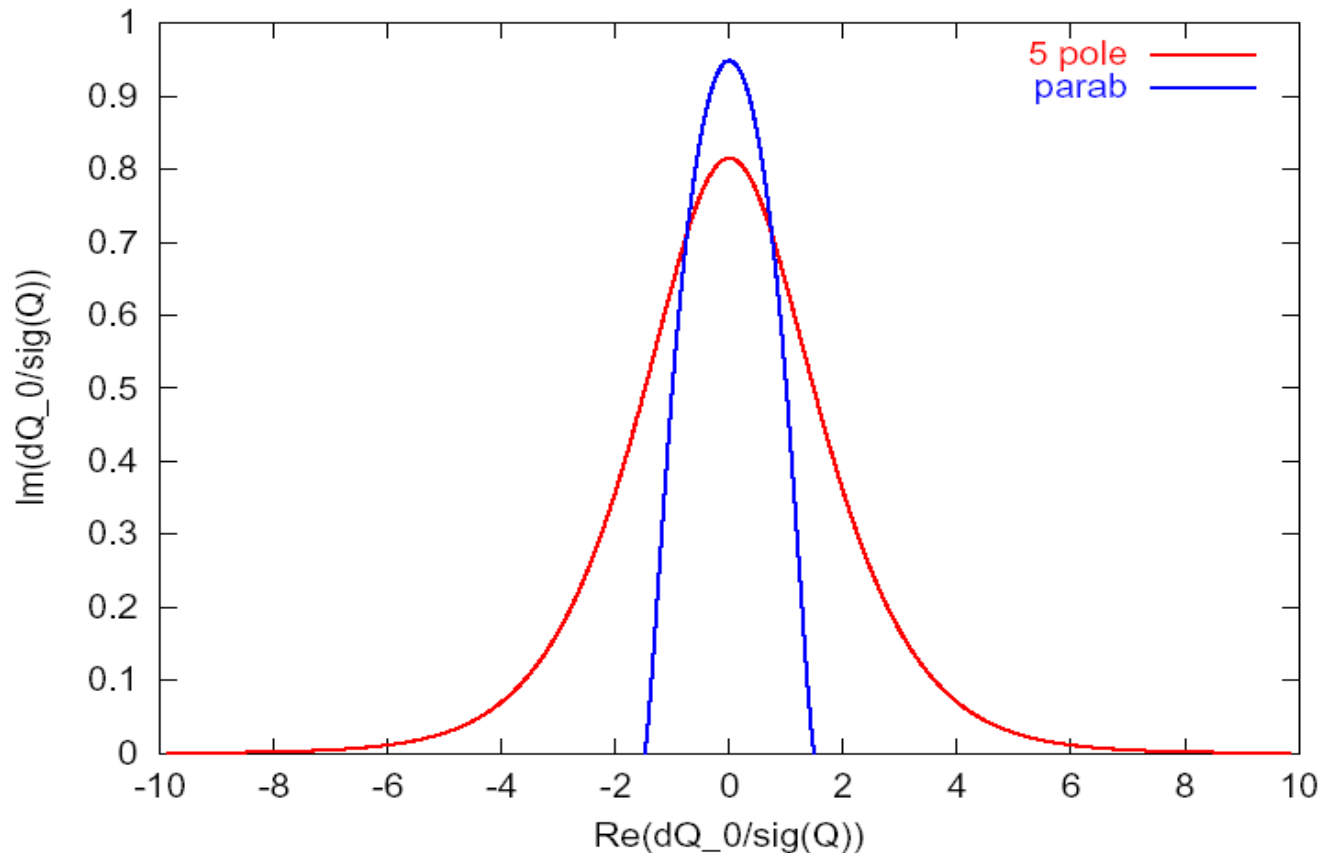
$$\delta Q_\beta = |\eta| Q_e \frac{\Delta p}{p} \Big|_{rms}$$

$$1 = \Delta Q_0 \int \frac{\rho_{norm}(v) dv}{\Delta Q + i0 - v \delta Q_\beta}$$

# Devil is in the details

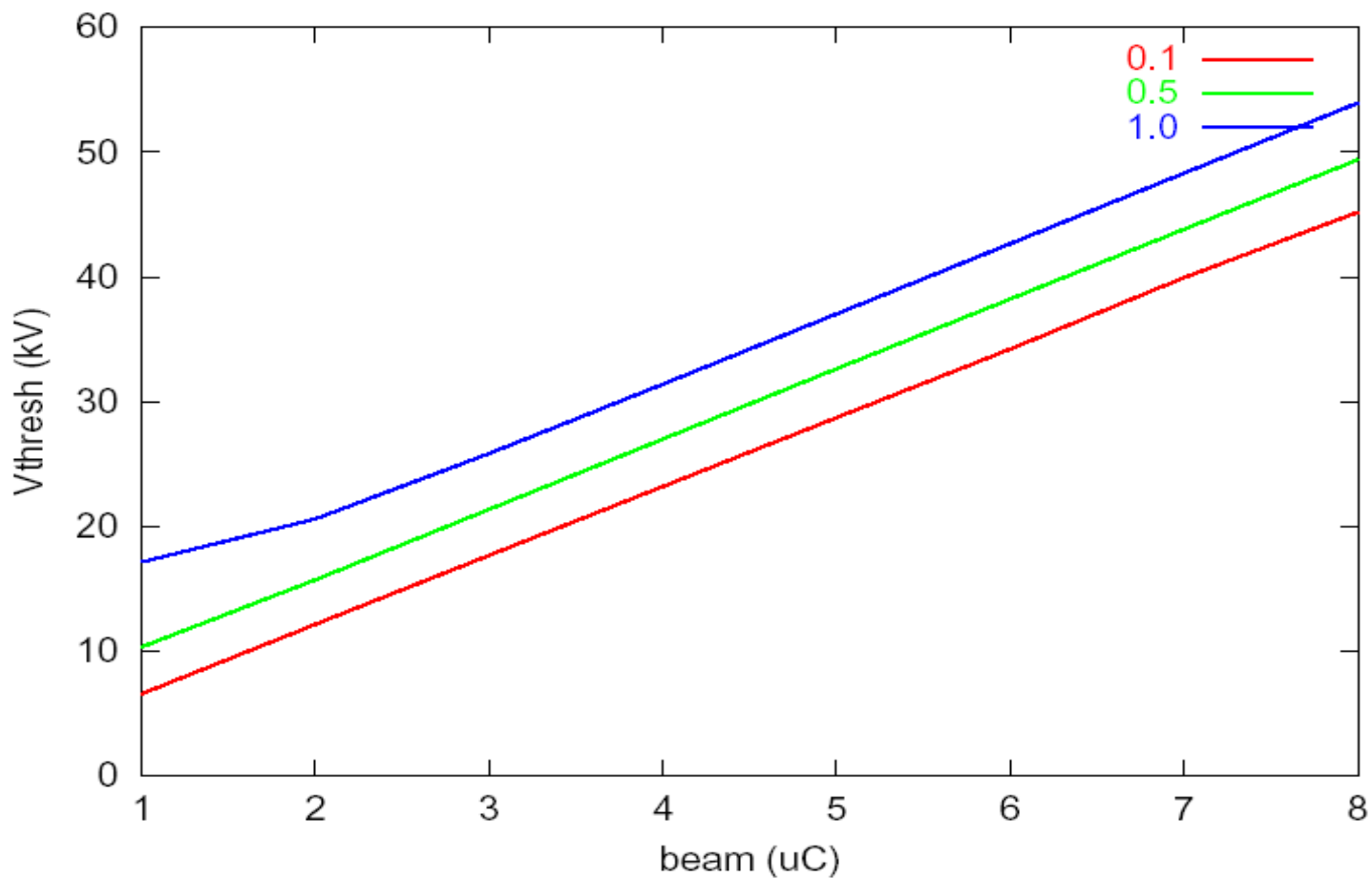
The momentum distribution used in the eigenmode analysis has long tails.  
This has a strong effect on the coasting beam dispersion relation.

Compare eigenmode distribution to parabolic with same RMS tune spread



# PSR is less stable with parabolic

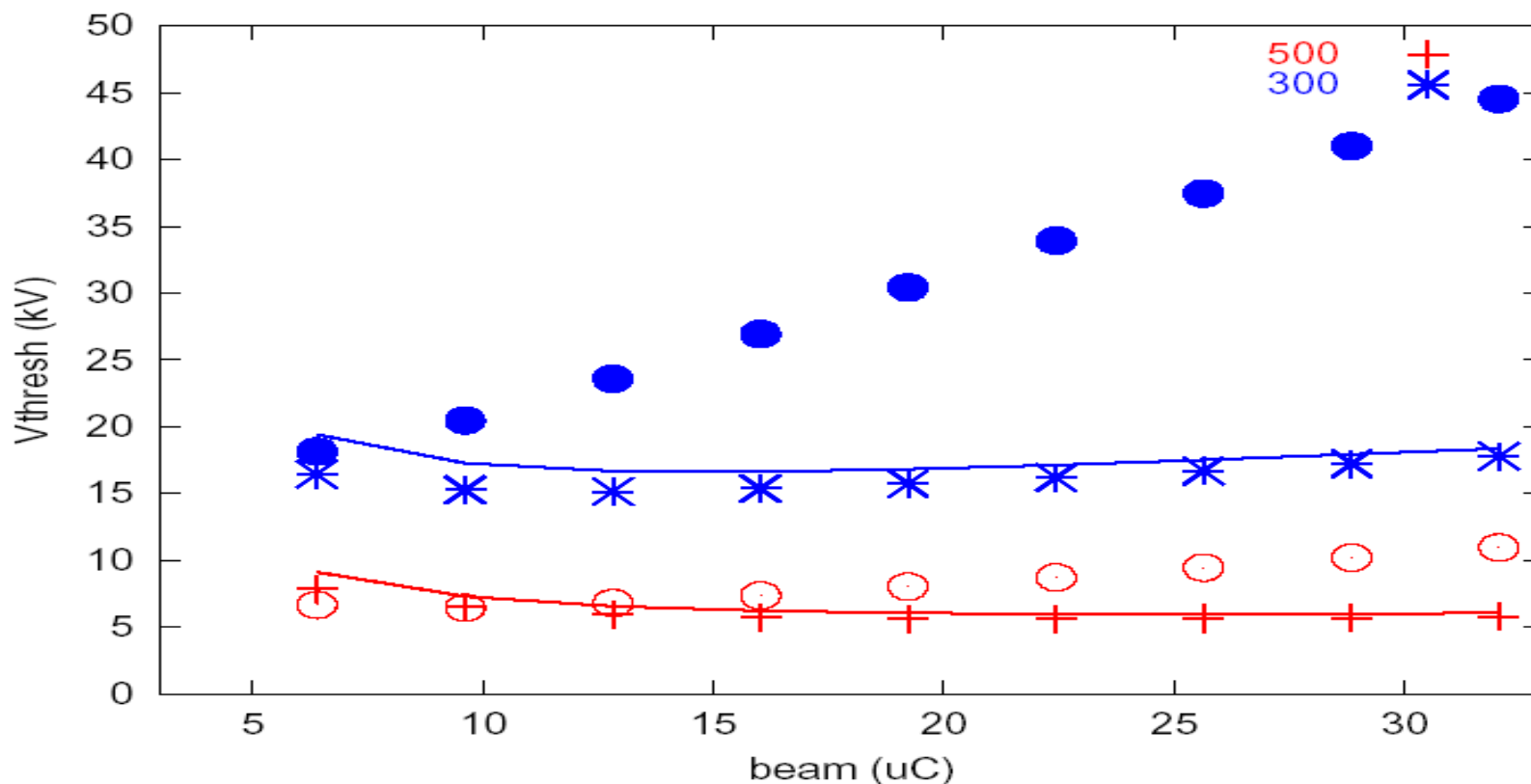
PSR with same inputs as previous page except using a parabolic momentum distribution



# SNS still looks good

2 nC/m, for nominal bunch equivalent square has 500 ns length (700 ns true)

Lines are coasting beam 5-pole, stars are eigenmode analysis, circles are coasting beam with parabolic momentum density



Electron cloud simulations and measurements in PSR are in fair agreement, around 1 nC/m of electrons survive gap

Simulations of e-p instability in PSR show threshold for 0.2 nC/m, the reason for the discrepancy is unclear

Simulations of 2MW SNS are stable for 1 nC/m

1 nC/m is for fairly large reflectivity and SEY.

Eigenmode analysis for SNS looks good with >2 nC/m

Scrubbing effects will improve things even more.

# Convergence test

$$\omega_e \tau_b = 5\pi$$

$$\Delta Q_0 = 0.10 + i0.05$$

$$\omega_0 \tau_b = 2\pi / 35$$

Needed many modes to see the  
beam breakup limit

Far fewer modes needed for  
reasonable threshold  
estimate

Coasting beam dispersion  
relation gave a good  
threshold estimate

