

Microscopic Phenomenological Model for the Secondary Emission Process

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- Motivation
- Definition of the model
- Fits to Cu and St. St. data
- Applications
 - time dependence of the electron-cloud power deposition in an LHC arc dipole
 - time dependence of the electron cloud dissipation in the PSR following extraction of the beam
- Conclusions



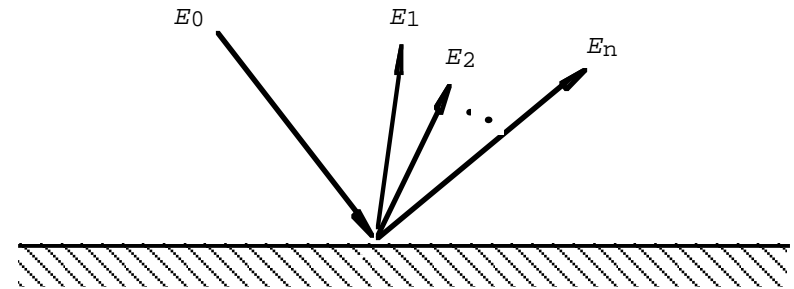
Our main focus:

- Understand in detail the effects of:
 - details of the secondary electron yield (SEY)
 - details of the secondary electron energy spectrum
- Simulations started in 1995 (with G. Lambertson):
 - PEP-II, LHC, APS, SPS, ...
 - main input: SEY (δ) and secondary energy spectrum ($d\delta/dE$)
 - obtained from lab measurements (R. Kirby, SLAC; N. Hilleret, CERN)
 - also books by Bruining, Redhead et al, and various articles in J. Appl. Phys. and other journals
- SEY model in the simulation originally developed in 1995
 - several refinements over past 1-1/2 years (M. Pivi)
 - some errors fixed (no major effect on results)



Details of the model:

- event-by-event simulation^(*)
 - event=one electron-wall collision
 - instantaneous generation of n secondaries (or absorption)
 - include E_0 and θ_0 dependence
- constraints for any single event:
 - $E_k \leq E_0$, $k=1,2,\dots,n$
 - $\sum_k E_k \leq E_0$
- constraints for the average over many events:
 - $\sum n P_n = \delta(E_0, \theta_0) = \text{input data}$
 - $d\delta/dE = \text{input data}$



(*) In our ECE simulation code, *all* macroparticles have the same charge (and mass)

Microscopic Probabilistic Model for the Simulation of Secondary Electron Emission*

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Abstract

We provide a detailed description of a model and its computational algorithm for the secondary electron emission process that has been used for some time now in the simulations of the electron cloud effect (ECE). We provide the numerical values for several parameters which have been obtained by fitting this model to laboratory measurements of the secondary yield and emitted energy spectrum.

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Basic macroscopic phenomenology

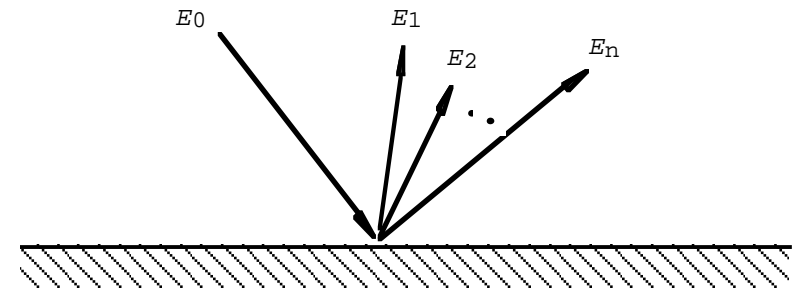
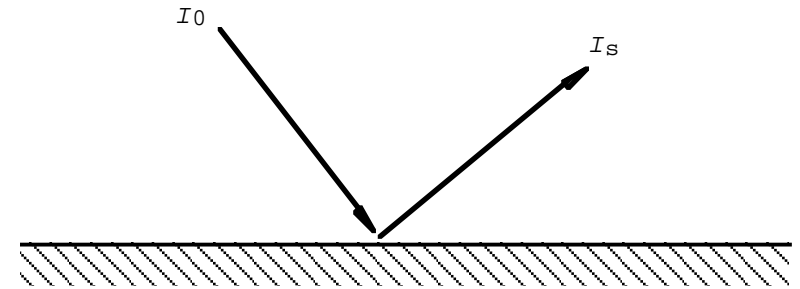
$$\text{SEY: } \delta = \frac{I_s}{I_0} \quad (\text{steady-state currents})$$

Cumulative energy spectrum:
apply a retarding voltage $E=eV$

$$S(E_0, E) = \frac{[I_s]_{\text{energy} \geq eV}}{I_0}$$

Retarding-voltage spectrum:

$$\left(\frac{d\delta}{dE} \right)_{\text{RV}} = - \frac{\partial S(E_0, E)}{\partial E}$$



What we need for an event-by-event simulation is:

$$\mathcal{P}_n = \frac{dP_n}{dE_1 d\Omega_1 dE_2 d\Omega_2 \cdots dE_n d\Omega_n}, \quad n = 1, 2, \dots$$



From \mathcal{P}_n obtain any macroscopic quantity:

$$P_n = \int (dE)_n (d\Omega)_n \mathcal{P}_n, \quad n \geq 1$$

$$\delta = \langle n \rangle = \sum_{n=1}^{\infty} n P_n$$

$$\left(\frac{d\delta}{dE} \right)_{RV} = \sum_{n=1}^{\infty} \int (dE)_n (d\Omega)_n \mathcal{P}_n \sum_{k=1}^n \delta(E_k - E)$$

We have constructed a phenomenological model for \mathcal{P}_n

- based as much as possible on data for δ and $d\delta/dE$
- not unique (use simplest assumptions whenever data is not available)
- many adjustable parameters, fixed by fitting δ and $d\delta/dE$ to data
- satisfies the single-event constraints
 - $E_k \leq E_0, k=1,2,\dots,n$
 - $\sum_k E_k \leq E_0$
- still evolving
- good agreement with data (by construction)
- some shortcomings



Ignore for now the emission-angle distribution (*ie.*, integrate over angles)

$$\frac{dP_n}{(dE)_n} = \theta(E_0 - E_1 - \dots - E_n) \prod_{k=1}^n f_n(E_k) \theta(E_k) \theta(E_0 - E_k)$$

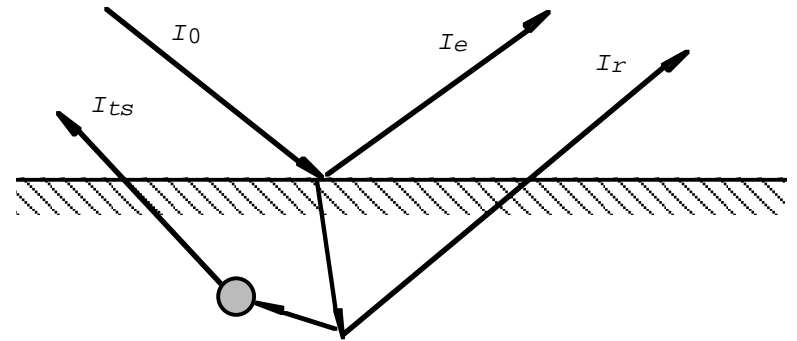
$$\int (dE)_n \frac{dP_n}{(dE)_n} = P_n$$

Three components of emitted electrons:

backscattered: $\delta_e = \frac{I_e}{I_0}$,

rediffused: $\delta_r = \frac{I_r}{I_0}$,

true secondaries: $\delta_{ts} = \frac{I_{ts}}{I_0}$



$$f_{1,e} = \theta(E)\theta(E_0 - E) \delta_e(E_0, \theta_0) \frac{2e^{-(E-E_0)^2/2\sigma_e^2}}{\sqrt{2\pi}\sigma_e \operatorname{erf}(E_0/\sqrt{2}\sigma_e)}$$

$$f_{1,r} = \theta(E)\theta(E_0 - E)\delta_r(E_0, \theta_0) \frac{(q+1)E^q}{E_0^{q+1}}$$

$$f_{n,ts} = \theta(E)F_n E^{p_n-1} e^{-E/\epsilon_n}$$

$$\delta_e(E_0, 0) = P_{1,e}(\infty) + (\hat{P}_{1,e} - P_{1,e}(\infty))e^{-(|E_0 - \hat{E}_e|/W)^p/p}$$

$$\delta_r(E_0, 0) = P_{1,r}(\infty) \left[1 - e^{-(E_0/E_r)^r} \right]$$

$$\delta_{ts}(E_0, \theta_0) = \hat{\delta}(\theta_0)D(E_0/\hat{E}(\theta_0)) \quad D(x) = \frac{s\mathbf{x}}{s-1+x^s}$$

$$P'_{n,ts} = \binom{M}{n} p^n (1-p)^{M-n}, \quad 0 \leq n \leq M \quad p = \langle n \rangle / M = \delta'_{ts} / M.$$

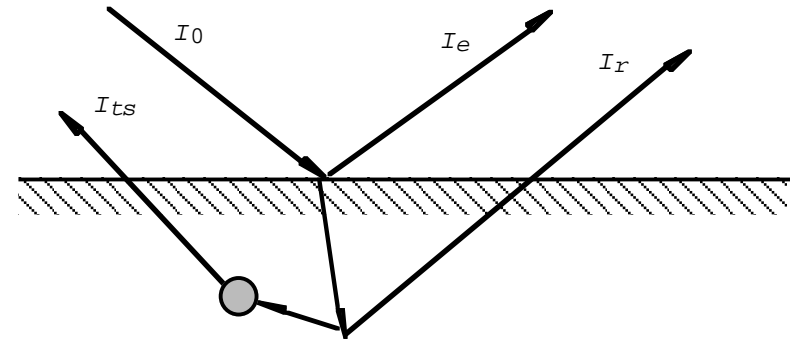
$$\delta'_{ts} = \frac{I_{ts}}{I_0 - I_e - I_r} = \frac{\delta_{ts}}{1 - \delta_e - \delta_r}$$

$$P_0 = 1 - \sum_{n=1}^{\infty} P_n$$



Exclusion assumption:

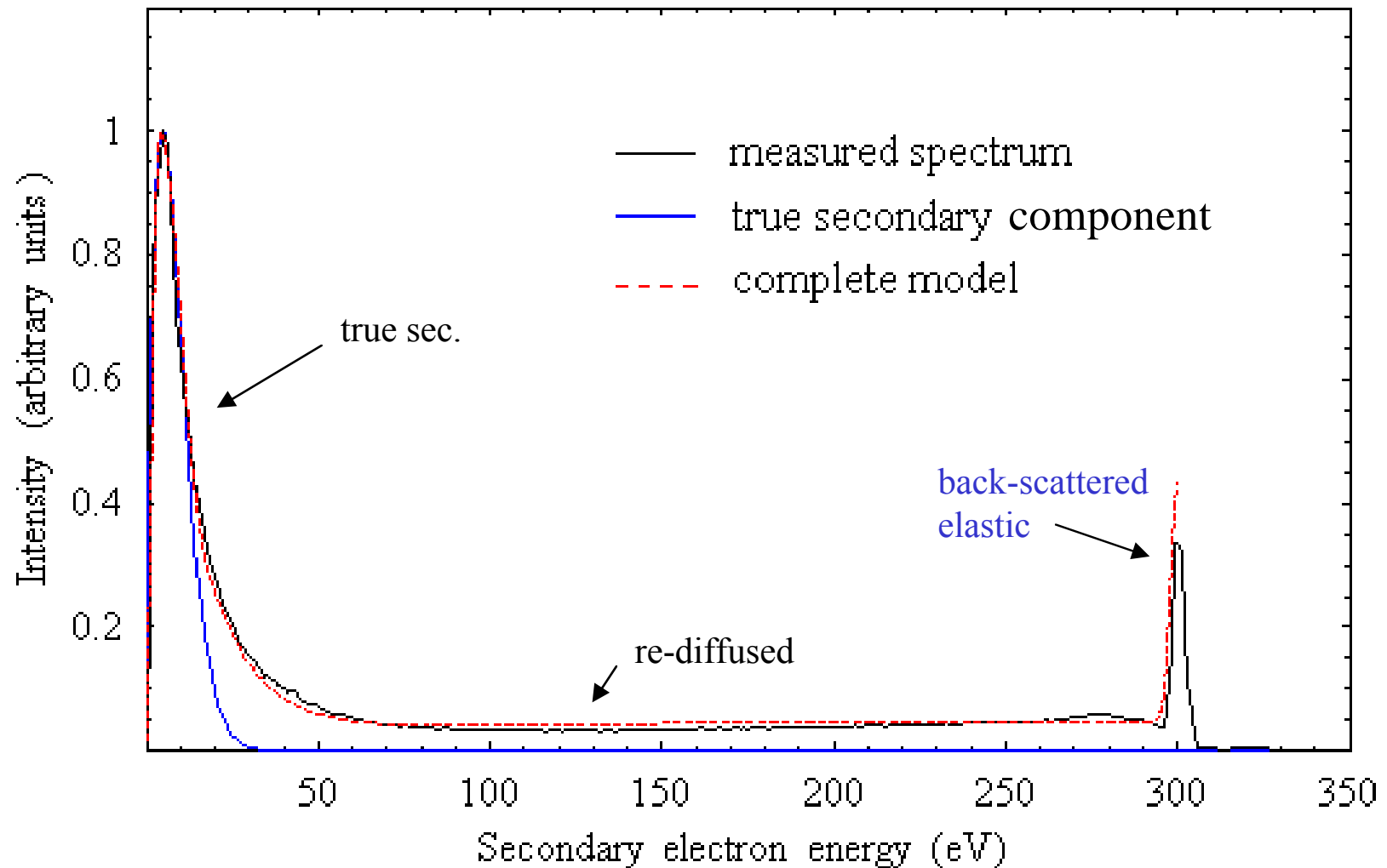
- if $n=1$, electron can be either backscattered, rediffused or true secondary
- if $n \geq 2$, electron can only be true secondary



This is a simplifying assumption:

- allows to easily extract the energy distribution functions $f_n(E)$ directly from the energy spectrum $d\delta/dE$
- no fundamental physical reason
- consistent with data on $d\delta/dE$ for $E_0 > 50$ eV or so, but not implied by it

Typical secondary emitted-electrons energy spectrum



Algorithm:

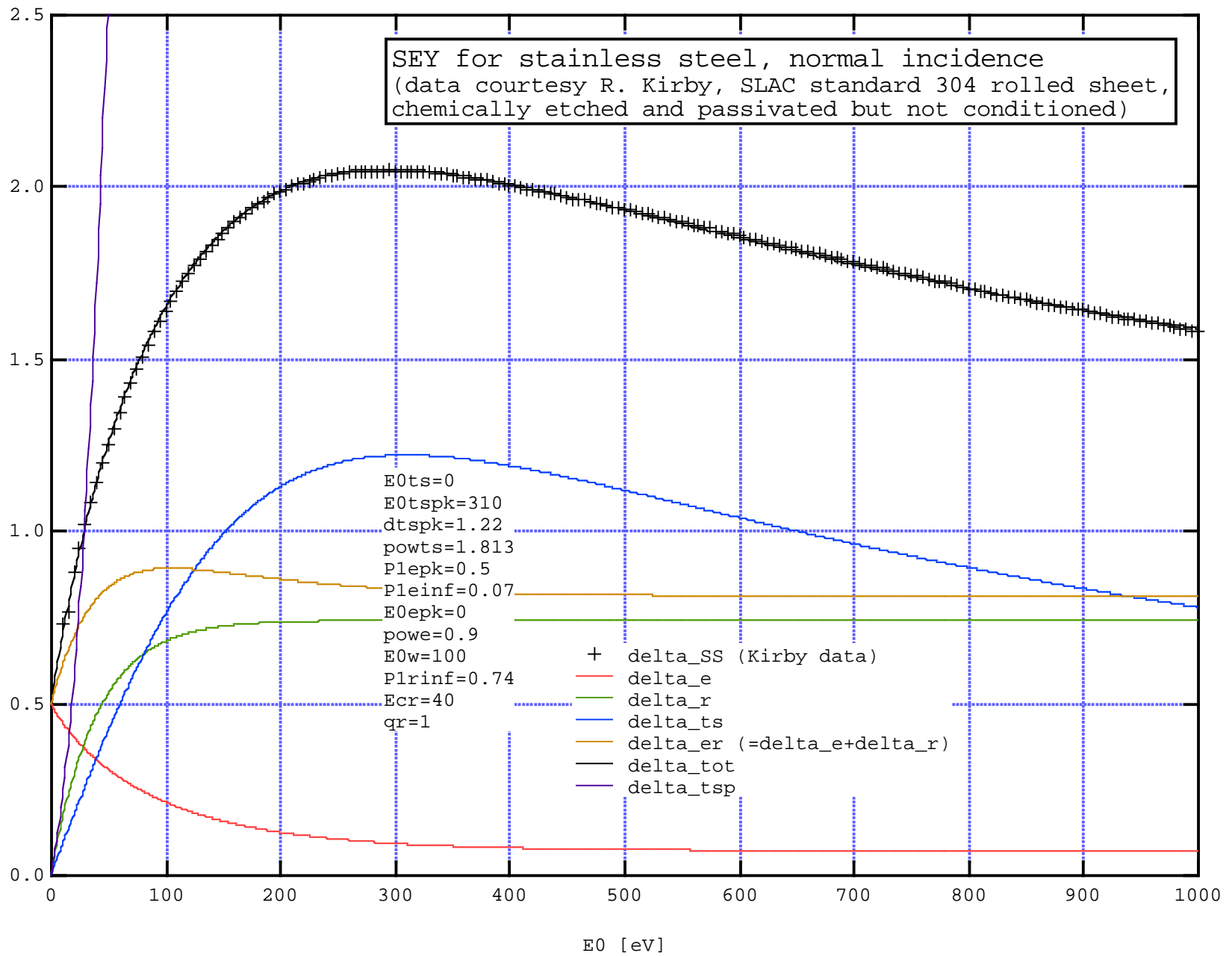
1. When an electron hits the wall, record E_0 and θ_0
2. Compute $\delta_e(E_0, \theta_0)$, $\delta_r(E_0, \theta_0)$ and $\delta_{ts}(E_0, \theta_0)$
3. Compute P_n , $n=1, 2, \dots, M$ (typically, $M=10$)
4. Generate a random integer $n \in [0, M]$ with probability distribution $\{P_n\}$
5. If $n=0$, electron got absorbed; continue with next incident electron
6. If $n=1$, generate its energy E with probability distribution $f_{1e}(E) + f_{1r}(E) + f_{1ts}(E)$ such that $E \leq E_0$
7. If $n \geq 2$, generate their energies E_k with probability distribution $f_{nts}(E_k)$ such that $\sum_{k=1}^n E_k \leq E_0$
8. Generate the emission angles (θ_k, ϕ_k) for $k=1, 2, \dots, n$
9. Continue with next incident electron

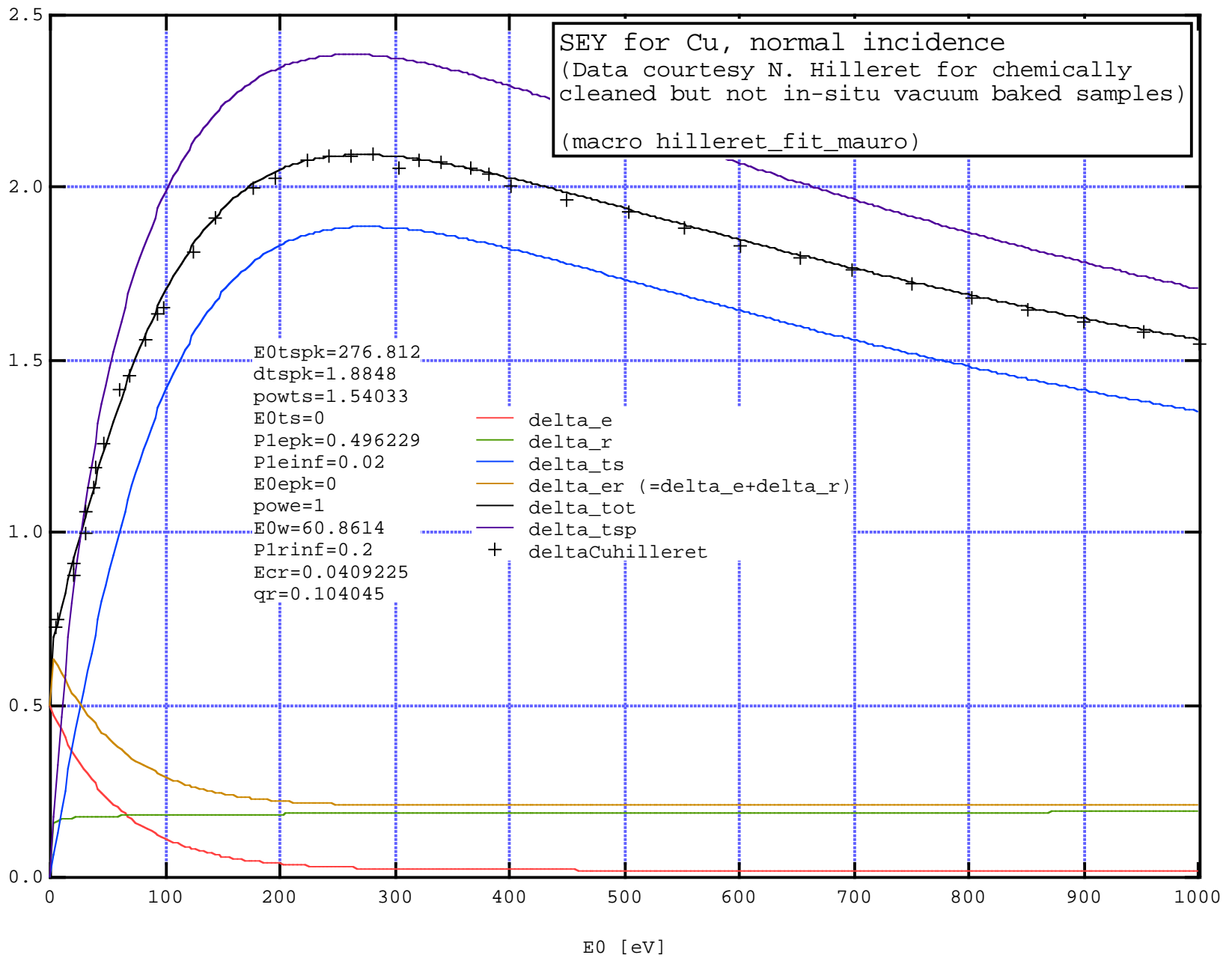


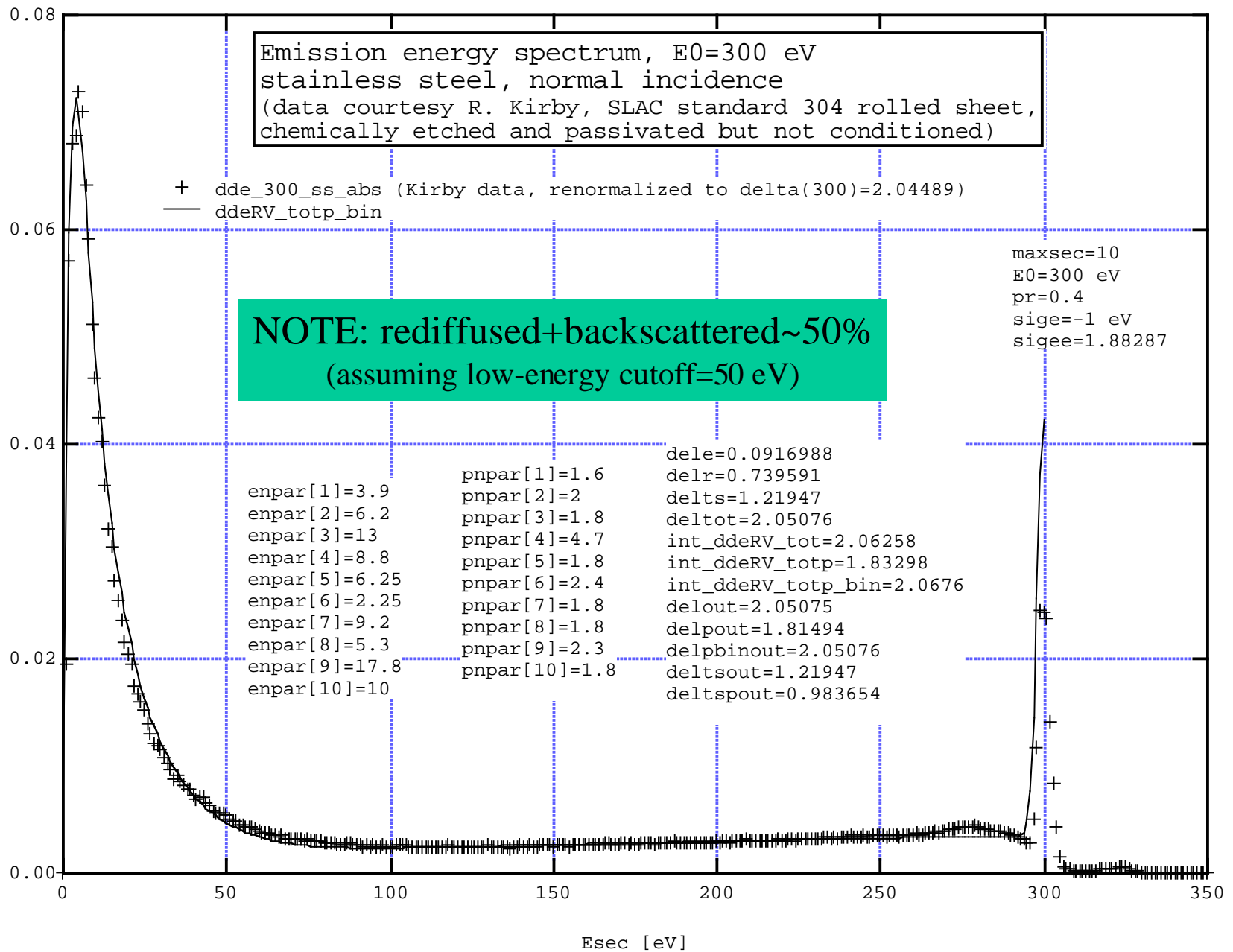
Fits to data:

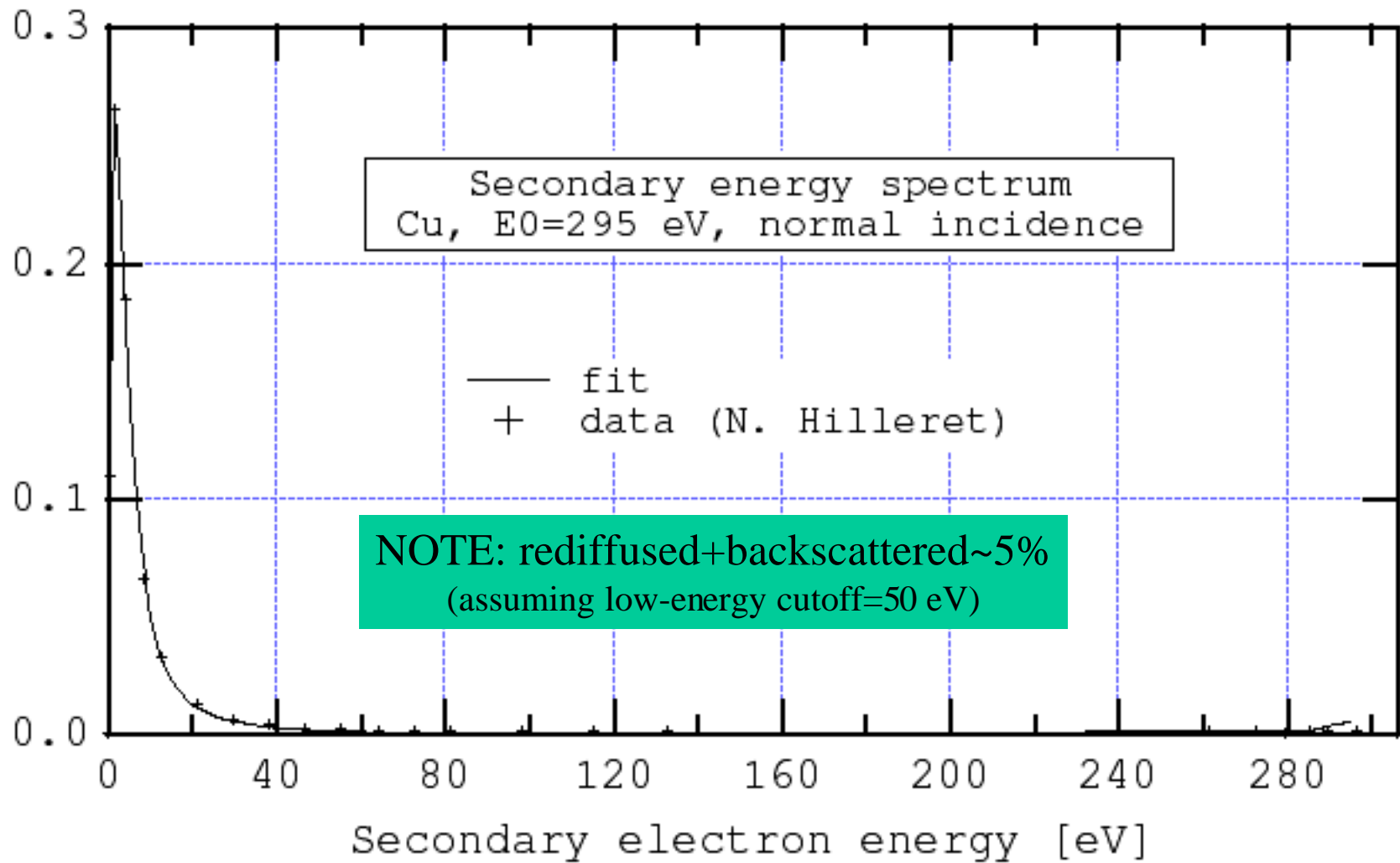
- Copper (data courtesy N. Hilleret, CERN)
 - sample: chemically cleaned but not vacuum-baked *in situ*
- Stainless steel (data courtesy R. Kirby, SLAC)
 - sample: SLAC standard 304 rolled sheet, chemically etched and passivated but not conditioned

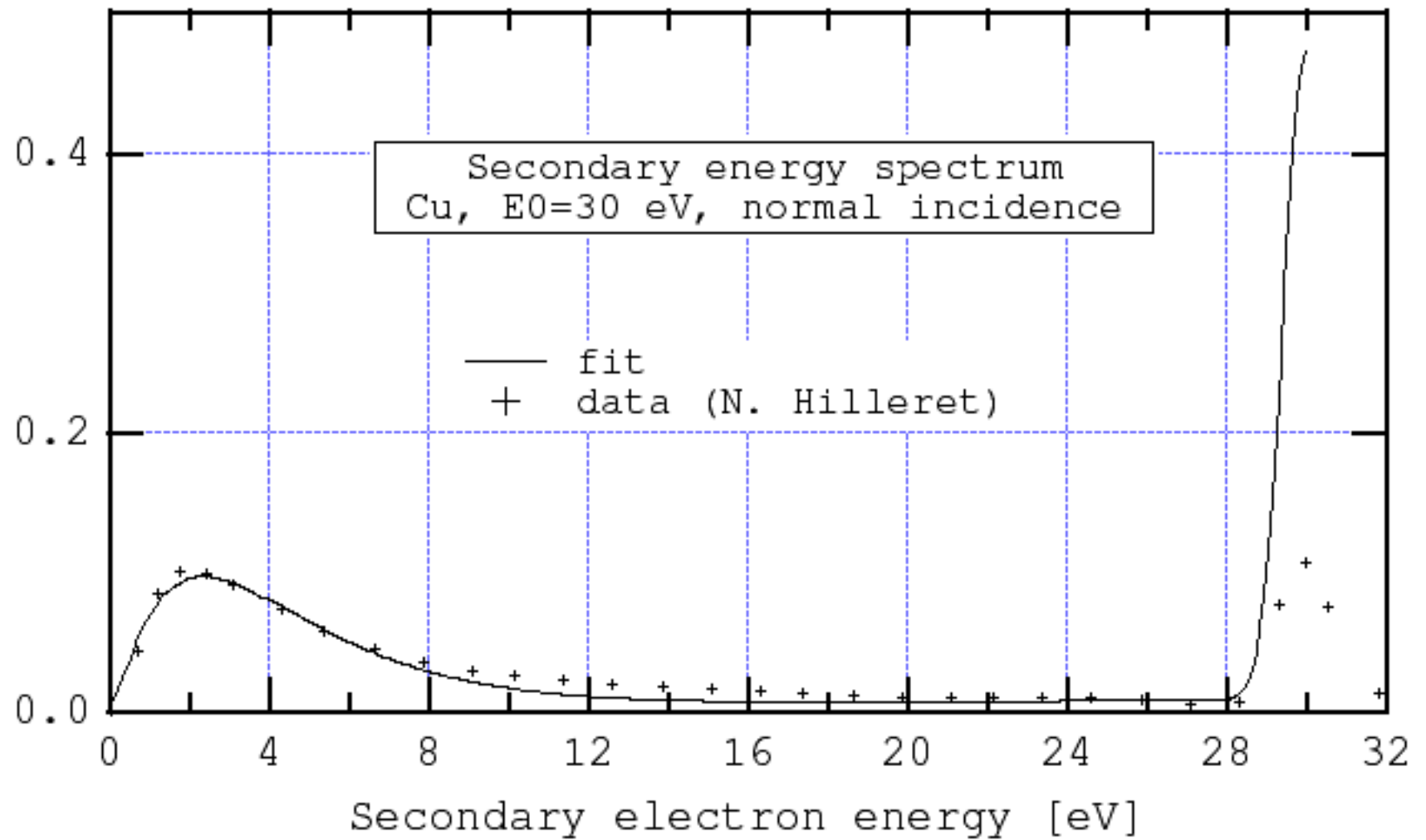


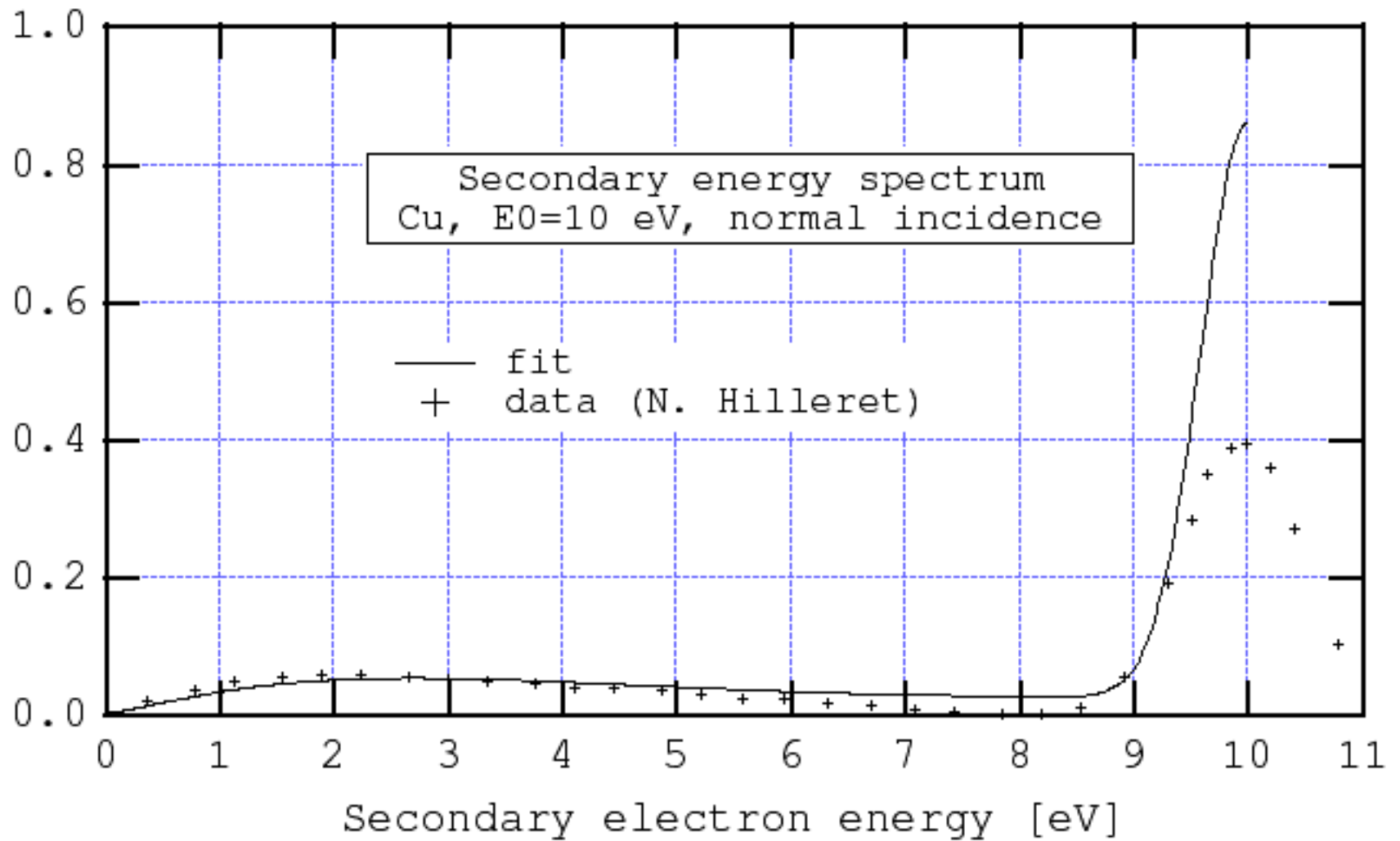












Current parameter values from fits to data

Table 1: Model parameters.

	Cu	SS	POSINST name
Emission angular spectrum (Sec. 2.3.1)			
α	1	1	pangsec
Elastically backscattered electrons (Sec. 3.2)			
$P_{1,e}(\infty)$	0.02	0.07	P1einf
$\hat{P}_{1,e}$	0.496	0.5	P1epk
\hat{E}_e [eV]	0	0	E0epk
W [eV]	60.86	100	E0w
p	1	0.9	powe
σ_e [eV]	2	1.9	sige
c_1	0.26	0.26	epar1
c_2	2	2	epar2
Rediffused electrons (Sec. 3.3)			
$P_{1,r}(\infty)$	0.2	0.74	P1rinf
E_r [eV]	0.041	40	Ecr
r	0.104	1	qr
q	0.5	0.4	pr
r_1	0.26	0.26	rpar1
r_2	2	2	rpar2
True secondary electrons (Sec. 3.4)			
$\hat{\delta}_{ts}$	1.8848	1.22	dtspk
\hat{E}_{ts} [eV]	276.8	310	E0tspk
s_0	1.54	1.813	powts
t_1	0.66	0.66	tpar1
t_2	0.8	0.8	tpar2
t_3	0.7	0.7	tpar3
t_4	1	1	tpar4
t_5	0	0	tpar5
t_6	0	0	tpar6
Total SEY[†]			
\hat{E}_t [eV]	271	292	
$\hat{\delta}_t$	2.1	2.05	dtotpk

[†] Note that $\hat{E}_t \simeq \hat{E}_{ts}$ and $\hat{\delta}_t \simeq \hat{\delta}_{ts} + P_{1,e}(\infty) + P_{1,r}(\infty)$ provided that $\hat{E}_{ts} \gg \hat{E}_e, E_r$.

Table 2: Further model parameters for the true secondary component.

	Cu	POSINST name
p_n	2.5, 3.3, 2.5, 2.5, 2.8, 1.3, 1.5, 1.5, 1.5, 1.5	pnpn(n)
c_n [eV]	1.5, 1.75, 1, 3.75, 8.5, 11.5, 2.5, 3, 2.5, 3	enpn(n)
	SS	POSINST name
p_n	1.6, 2, 1.8, 4.7, 1.8, 2.4, 1.8, 1.8, 2.3, 1.8	pnpn(n)
c_n [eV]	3.9, 6.2, 13, 8.8, 6.25, 2.25, 9.2, 5.3, 17.8, 10	enpn(n)



Q: is the electron emitted spectrum Maxwellian?

A: only approximately.

definition of Maxwellian spectrum:

$$\begin{aligned}\frac{dN}{d^3\mathbf{p}} &\propto \exp(-E/kT), & E &= \frac{\mathbf{p}^2}{2m_e} \\ \Rightarrow \frac{dN}{dE} &\propto E^{1/2} \exp(-E/kT) \equiv E^{p_n-1} \exp(-E/\varepsilon_n) \\ \Rightarrow p_n &= 3/2\end{aligned}$$

Fits to data, however, imply $p_n \sim 1.8-5$, depending on n and material

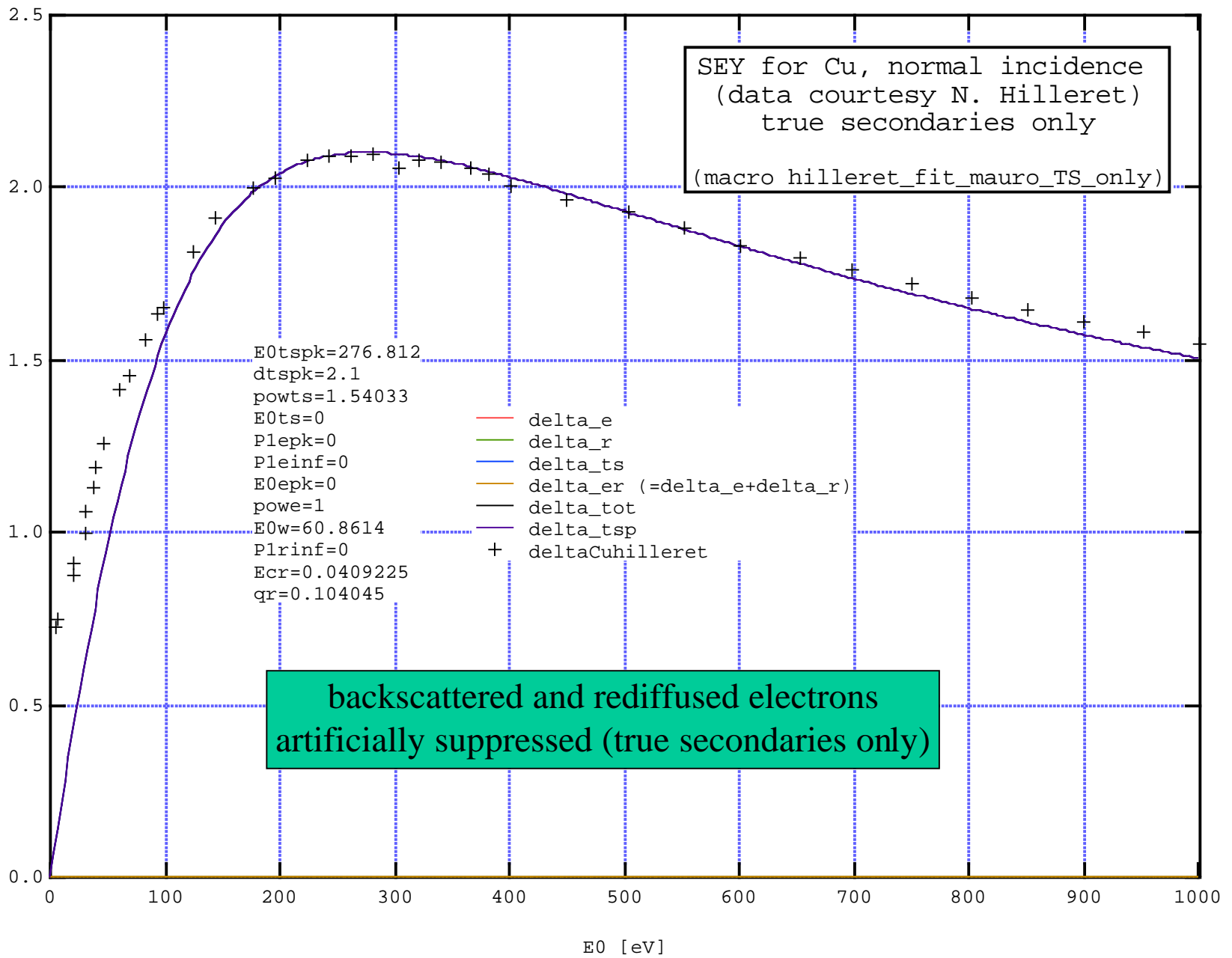


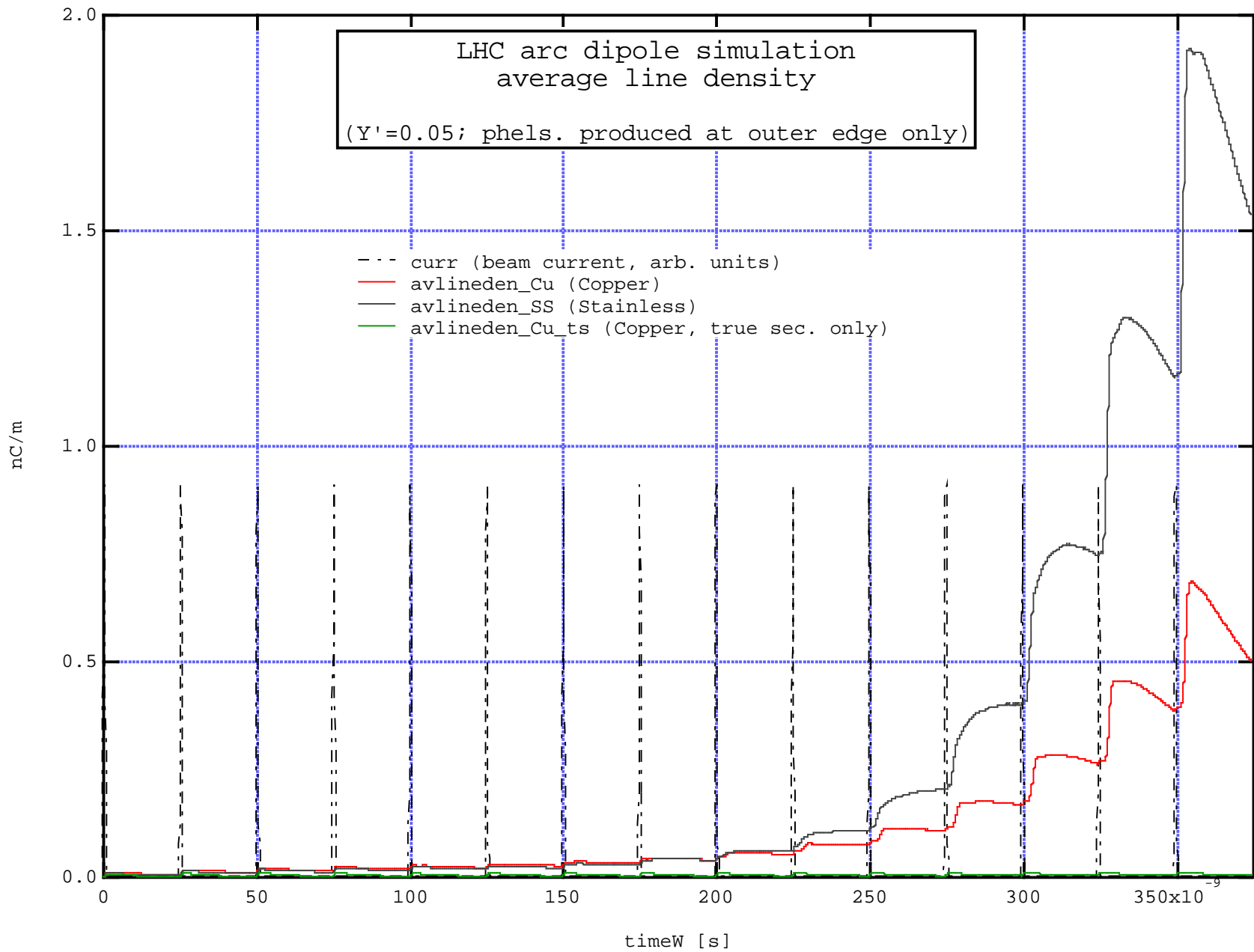
Applications:

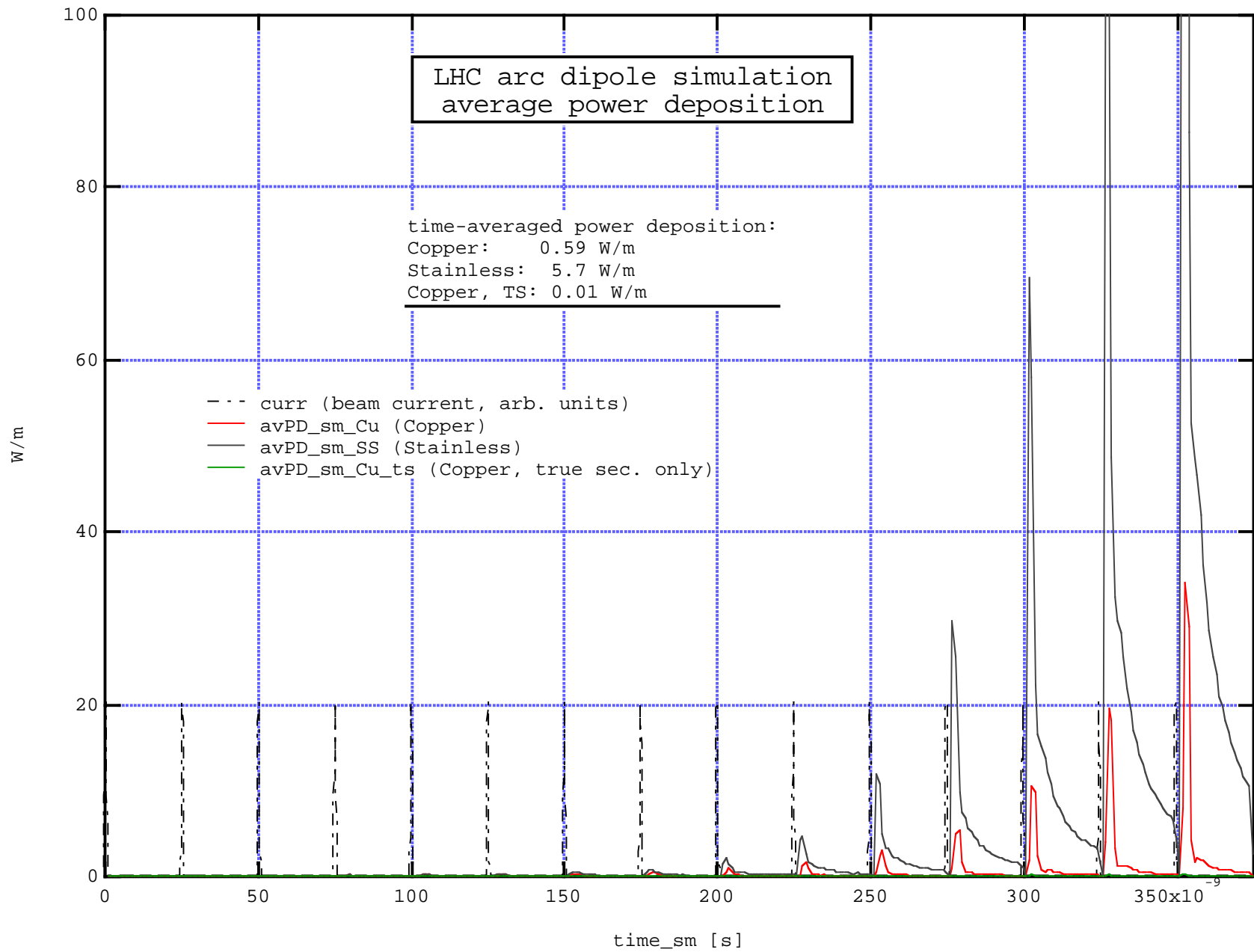
- (1) time dependence of the electron-cloud power deposition in LHC arc dipole:
simulate injection of a 15-bunch train into an empty chamber;
 - sensitivity to relative ratio of backscattered, rediffused and true secondaries
 - sensitivity to $\delta(0)$

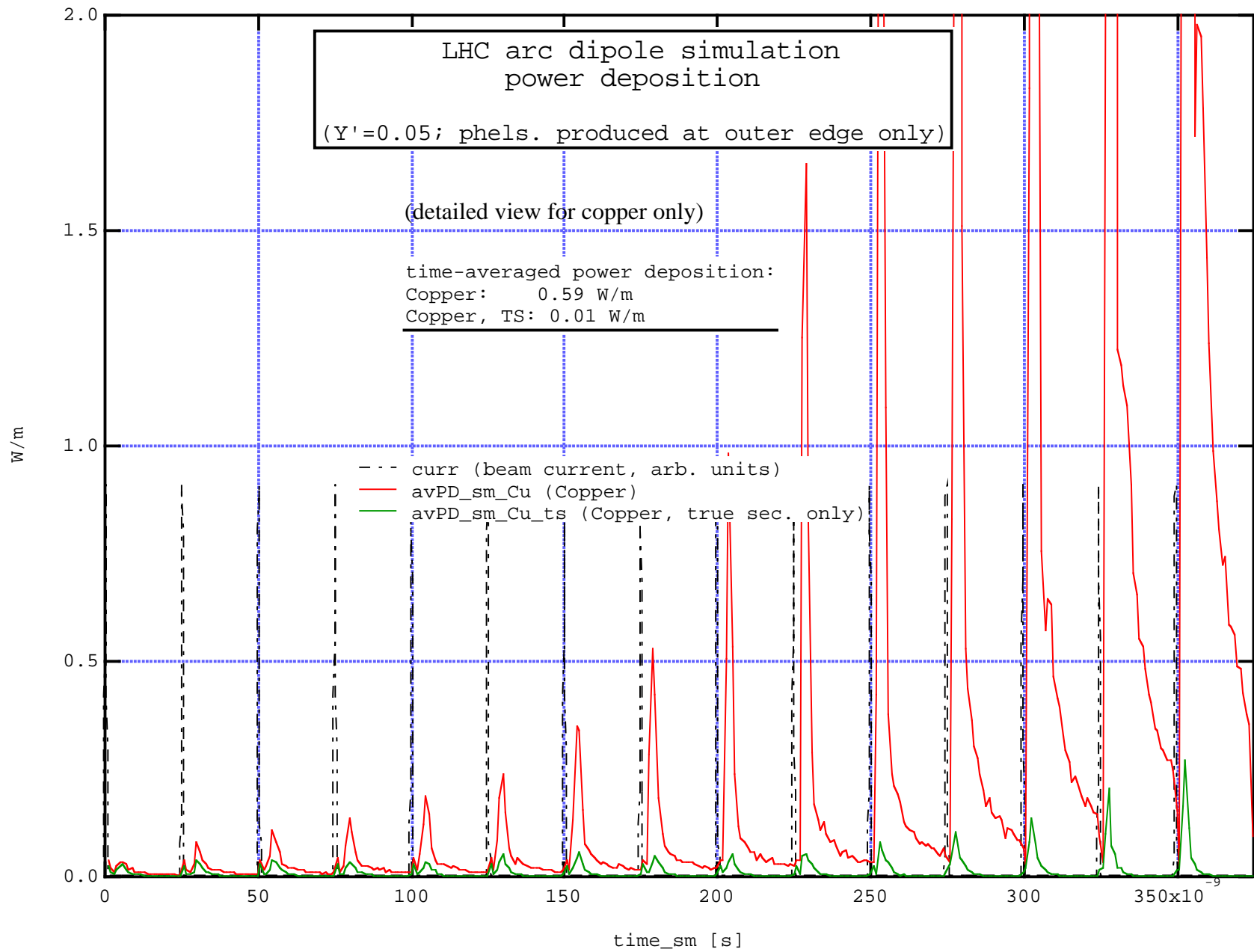
- (2) time dependence of the electron cloud dissipation following extraction of the PSR beam

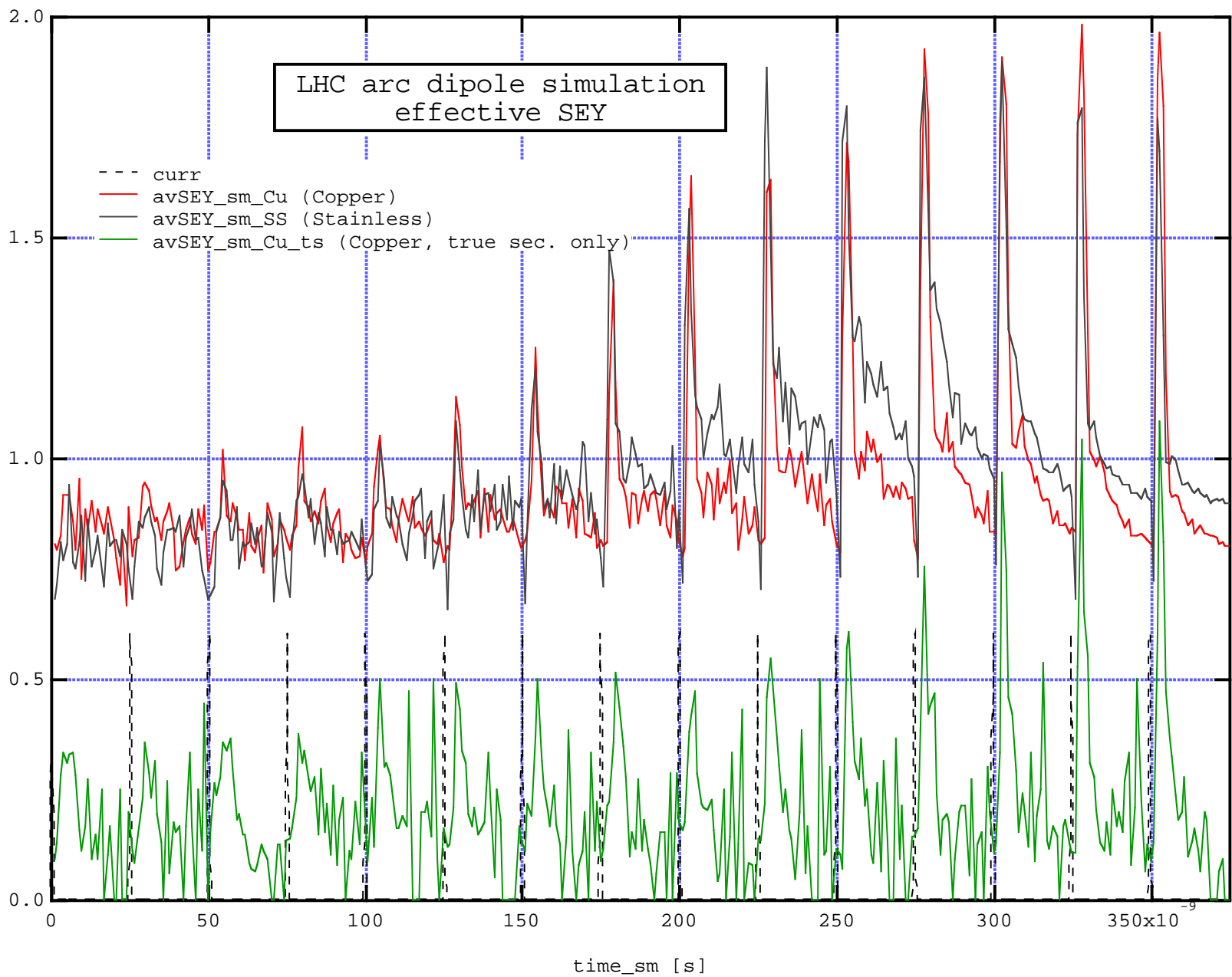


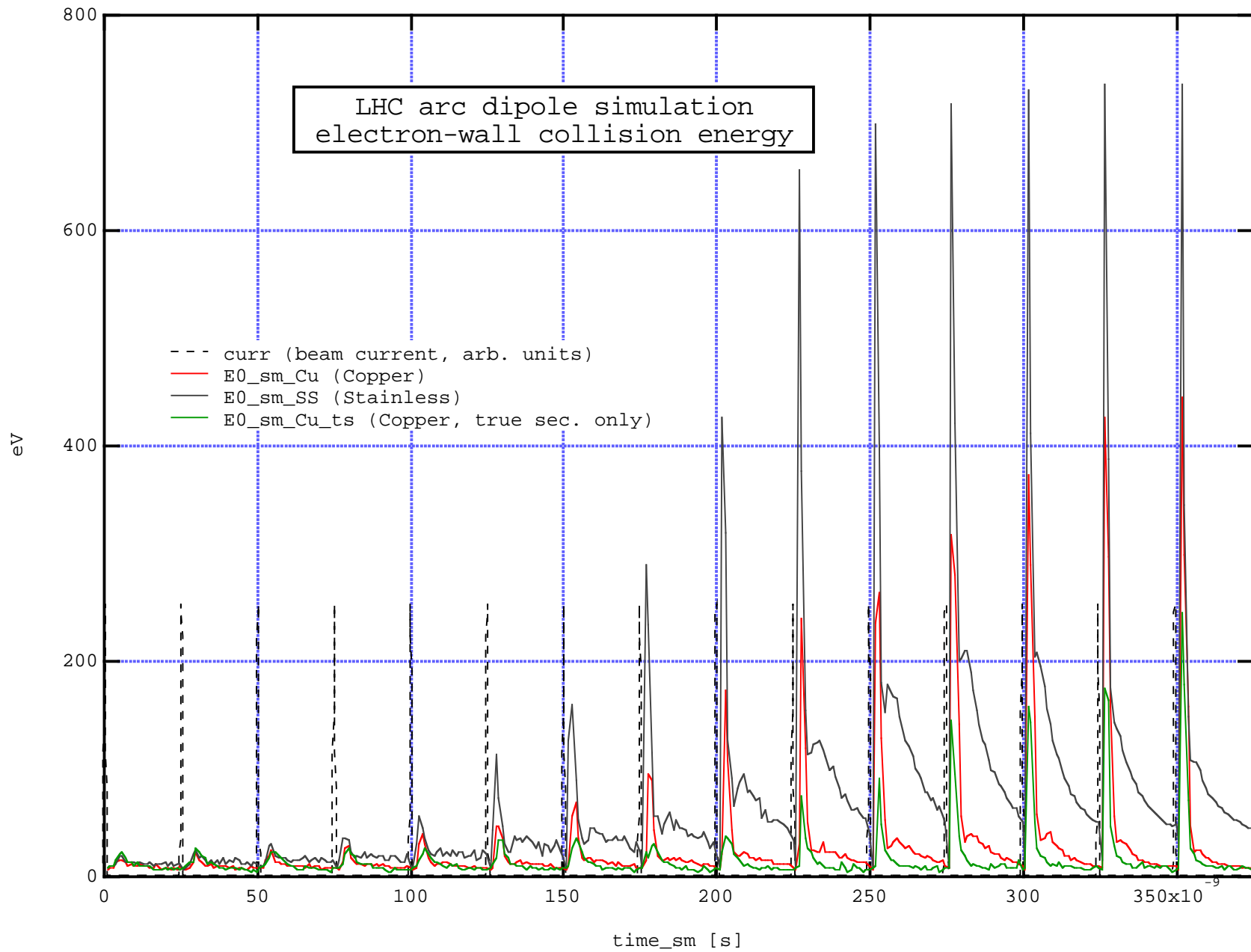


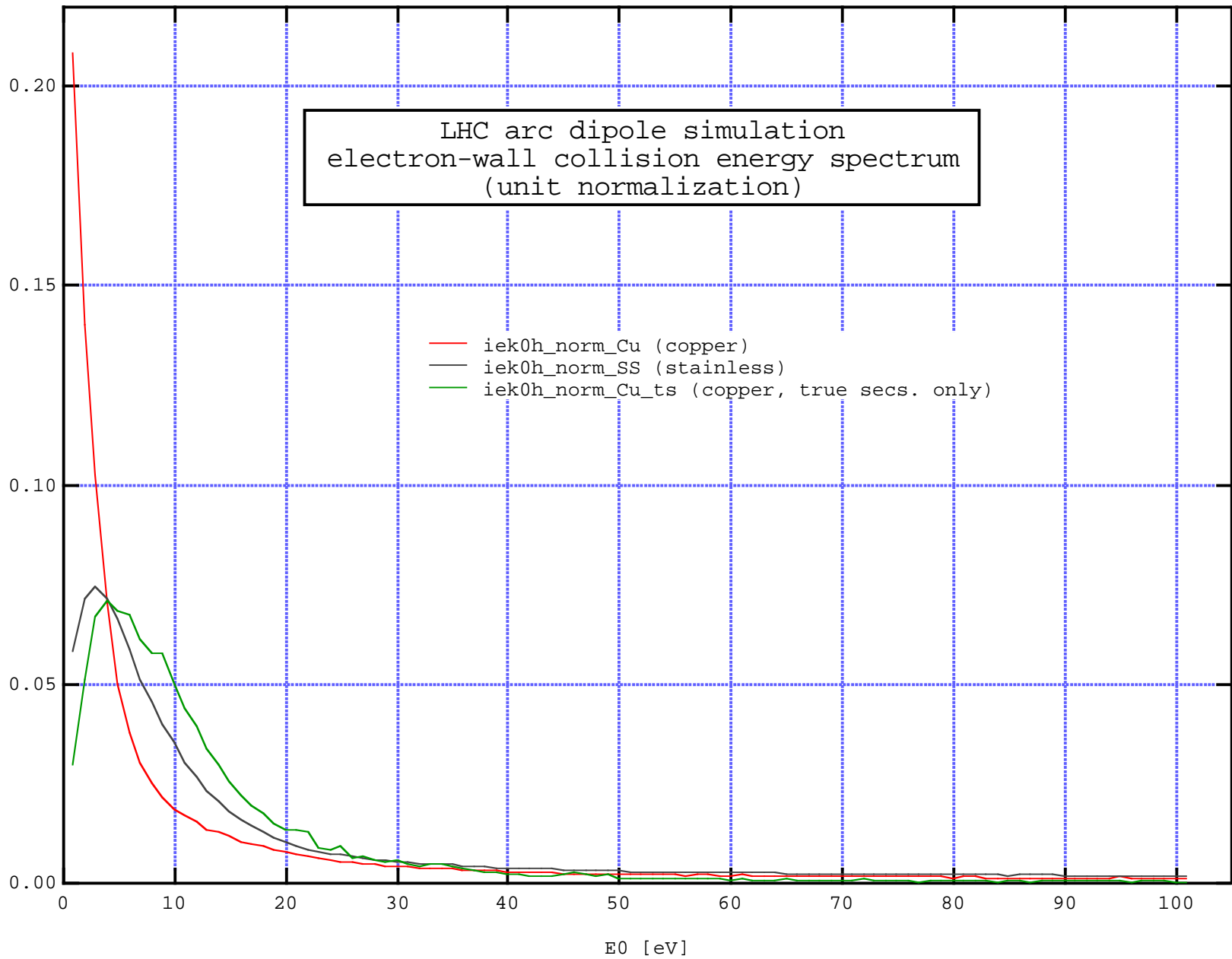




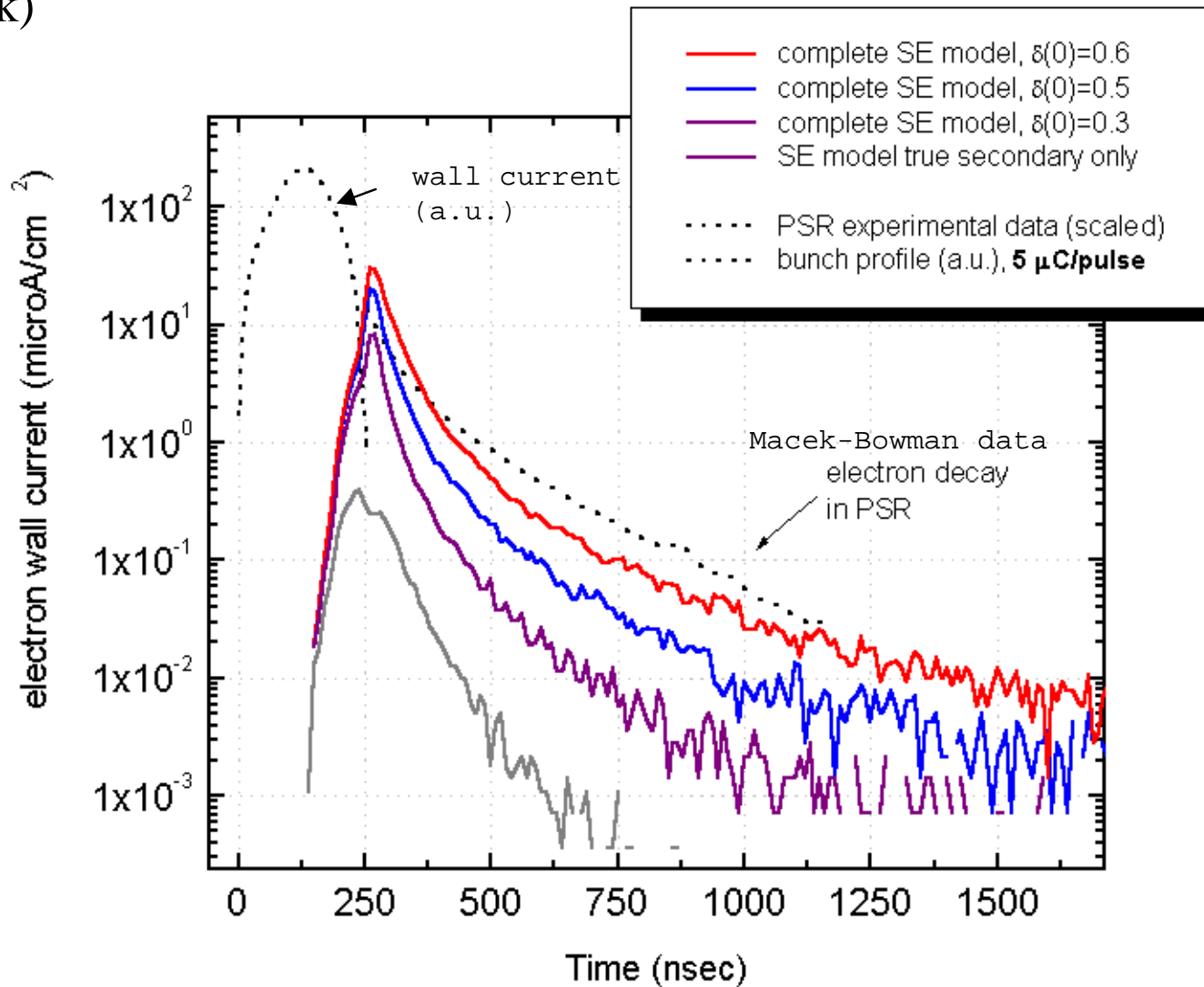








Dissipation of the electron cloud upon extraction of the PSR beam (data courtesy R. Macek)



Conclusions:

- constructed phenomenological model for secondary emission process
- consistent with data on $\delta(E_0)$ and $d\delta/dE$ (by construction)
- probabilities P_n calculated self-consistently (ie., $\delta(E_0)$ and $d\delta/dE$ are reproduced from the P_n 's)
- still evolving; not enough detailed data to pin down all parameters

- Applications:
 - power deposition in LHC dipoles exhibits a sensitivity to details of the secondary emission spectrum (relative ratio of backscattered, rediffused and true secondaries)
 - however, this sensitivity may be less pronounced when in steady state
 - also sensitivity to $\delta(0)$ (not new)
 - for PSR, infer $\delta_{\text{eff}}(0)=0.5-0.6$ from measured data on dissipation of the electron cloud following extraction



Sample comparison of simulation vs. measurements at the APS

